

Lattice approximation to the dynamical Φ_3^4 model *

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Abstract

We study the lattice approximations to the dynamical Φ_3^4 model by paracontrolled distributions proposed in [GIP13]. We prove that the solutions to the lattice systems converge to the solution to the dynamical Φ_3^4 model in probability, locally in time. The dynamical Φ_3^4 model is not well defined in the classical sense. Renormalisation has to be performed in order to define the non-linear term. Formally, this renormalisation corresponds to adding an infinite mass term to the equation which leads to adding a drift term in the lattice systems.

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1 Introduction

Recall that the usual continuum Euclidean Φ_d^4 -quantum field theory is heuristically described by the following probability measure:

$$N^{-1} \Pi_{x \in \mathbb{T}^d} d\phi(x) \exp \left(- \int_{\mathbb{T}^d} (|\nabla \phi(x)|^2 + \phi^2(x) + \phi^4(x)) dx \right), \quad (1.1)$$

where N is the normalization constant and ϕ is the real-valued field and \mathbb{T}^d is the d -dimensional torus. There have been many approaches to the problem of giving a meaning to the above heuristic measure for $d = 2$ and $d = 3$ (see Refs. [GRS75], [GJ87] and references therein). In [PW81] Parisi and Wu proposed a program for Euclidean quantum field theory of getting Gibbs states of classical statistical mechanics as limiting distributions of stochastic processes, especially as solutions to non-linear stochastic differential equations. Then one can use the

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stochastic differential equations to study the properties of the Gibbs states. This procedure is called stochastic field quantization (see [JLM85]). The Φ_d^4 model is the simplest non-trivial Euclidean quantum field (see [GJ87] and the reference therein). The issue of the stochastic quantization of the Φ_d^4 model is to solve the following equation:

$$d\Phi = (\Delta\Phi - \Phi^3)dt + dW(t) \quad \Phi(0) = \Phi_0. \quad (1.2)$$

where W is a cylindrical Wiener process on $L^2(\mathbb{T}^d)$. The solution Φ is also called dynamical Φ_d^4 model.

In two spatial dimensions, the dynamical Φ_2^4 model was previously treated in [AR91], [DD03] and [MW15]. In three spatial dimensions this equation (1.1) is ill-posed and the main difficulty in this case is that W and hence the solutions are so singular that the non-linear term is not well-defined in the classical sense. It was a long-standing open problem to give a meaning to equation (1.2) in the three dimensional case. A breakthrough result was achieved recently by Martin Hairer in [Hai14], where he introduced a theory of regularity structures and gave a meaning to equation (1.2) successfully. Also by using the paracontrolled distributions proposed by Gubinelli, Imkeller and Perkowski in [GIP13] existence and uniqueness of local solutions to (1.2) has been obtained in [CC13]. Recently, these two approaches have been successful in giving a meaning to a lot of ill-posed stochastic PDEs like the Kardar-Parisi-Zhang (KPZ) equation ([KPZ86], [BG97], [Hai13]), the dynamical Φ_3^4 model ([Hai14], [CC13]), the Navier-Stokes equation driven by space-time white noise ([ZZ14a], [ZZ14b]), the dynamical sine-Gordon equation ([HS14]) and so on (see [HP14] for more interesting examples). From a philosophical perspective, the theory of regularity structures and the paracontrolled distributions are inspired by the theory of controlled rough paths [Lyo98], [Gub04]. The main difference is that the regularity structure theory considers the problem locally, while the paracontrolled distribution method is a global approach using Fourier analysis. In [Kup14] the author also use renormalization group techniques to make sense of the dynamical Φ_3^4 model. .

The lattice approximation is an important tool in constructing and studying the continuum Φ_3^4 field (see [P75, P77, ABZ04]). It also preserves Osterwalder-Schrader positivity and also the ferromagnetic nature of the measure (see [GJ87] and the references therein). Let us set $\Lambda_\varepsilon := \{\varepsilon x \in \mathbb{T}^3, x \in \mathbb{Z}^3\}$. Heuristically, the quantities $\int |\nabla \phi(x)|^2 dx$, $\int \phi^2(x) dx$, and $\int \phi^4(x) dx$ can be approximated by $\varepsilon \sum_{|x-y|=\varepsilon, x, y \in \Lambda_\varepsilon} (\phi(x) - \phi(y))^2$, $\varepsilon^3 \sum_{x \in \Lambda_\varepsilon} \phi(x)^2$ and $\varepsilon^3 \sum_{x \in \Lambda_\varepsilon} \phi(x)^4$, respectively, as ε tends to zero. Thus heuristically (1.1) can be approximated by the following heuristic probability measure μ_ε :

$$N_\varepsilon^{-1} \Pi_{x \in \Lambda_\varepsilon} d\phi_x \exp \left(2\varepsilon \sum_{|x-y|=\varepsilon, x, y \in \Lambda_\varepsilon} \phi(x)\phi(y) - (\varepsilon^3 + 12\varepsilon) \sum_{x \in \Lambda_\varepsilon} \phi^2(x) - \varepsilon^3 \sum_{x \in \Lambda_\varepsilon} \phi^4(x) \right), \quad (1.3)$$

where N_ε is the normalization constant. (1.3) is still just a heuristic expression, but it is indeed not hard to give a rigorous sense to it (see [GJ87] and the references therein). We call this the lattice Φ_3^4 -field measure. From μ_ε by deriving suitable bounds on its moments and choosing subsequences if necessary one gets limit measures by weak convergence. These are then the continuum Φ_3^4 -field measures.

The following stochastic PDEs on Λ_ε are the stochastic quantizations associated with the

lattice Φ_3^4 -field measure:

$$\begin{aligned} d\Phi^\varepsilon(t, x) = & (\Delta_\varepsilon \Phi^\varepsilon(t, x) - (\Phi^\varepsilon)^3(t, x) + (3C_0^\varepsilon - 9C_1^\varepsilon)\Phi^\varepsilon(t, x))dt \\ & + \varepsilon^{-3/2}dW_\varepsilon(t, x) \\ \Phi^\varepsilon(0) = & \Phi_0^\varepsilon. \end{aligned} \tag{1.4}$$

Here $W_\varepsilon(t) = \{W(t, x)\}_{x \in \Lambda_\varepsilon}$ is a family of independent Brownian motions and C_0^ε and C_1^ε are defined as below. For $x \in \Lambda_\varepsilon$ define

$$\Delta_\varepsilon f(x) := \varepsilon^{-2} \sum_{y \in \Lambda_\varepsilon, y \sim x} (f(y) - f(x)),$$

and the nearest neighbor relation $x \sim y$ is to be understood with periodic boundary conditions on Λ_ε . We emphasize that to make sense of (1.2) we need to renormalise some ill-defined terms in (1.2). This is done by adding the renormalisation terms $C_0^\varepsilon \Phi^\varepsilon$ and $C_1^\varepsilon \Phi^\varepsilon$ in the approximating equation (1.4). In this paper we prove that the dynamical lattice approximation converge to the dynamical Φ_3^4 model. This problem is also related to the convergence of a rescaled discrete spin system to the solution to the dynamical Φ_3^4 model (see [MW14] for the dynamical Φ_2^4 model).

In one dimensional case, approximations to general stochastic partial differential equations driven by space-time white noise have been very well studied (see [Gy98, Gy99, DG01, HMW14] and the reference therein). In [GP15] the authors study the Sasamoto-Spohn type discretizations of the conservative stochastic Burgers equation. In three dimensional case we also study discrete approximations to N-S equations (see [ZZ14a]).

In this paper we use the paracontrolled distribution method to prove that the solutions to the lattice approximating equation converge to the solution of the dynamical Φ_3^4 model. The theory of paracontrolled distributions combines the idea of Gubinelli's controlled rough path [Gub04] and Bony's paraproduct [Bon84], which is defined as follows: Let $\Delta_j f$ be the j th Littlewood-Paley block of a (Scharwtz) distribution f . Define for distributions f and g

$$\pi_{<}(f, g) = \pi_{>}(g, f) = \sum_{j \geq -1} \sum_{i < j-1} \Delta_i f \Delta_j g, \quad \pi_0(f, g) = \sum_{|i-j| \leq 1} \Delta_i f \Delta_j g.$$

Formally $fg = \pi_{<}(f, g) + \pi_0(f, g) + \pi_{>}(f, g)$. Observing that if f is regular $\pi_{<}(f, g)$ behaves like g and is the only term in the Bony's paraproduct not improving the regularities, the authors in [GIP13] consider a paracontrolled ansatz of the type

$$u = \pi_{<}(u', g) + u^\sharp,$$

where $\pi_{<}(u', g)$ represents the "bad-part" of the solution, u' is some suitable function and g is some functional of the Gaussian field and u^\sharp is regular enough to define the multiplication required. Then to make sense of the product of uf we only need to define gf .

Using the paracontrolled distribution method, to perform the lattice approximation of the dynamical Φ_3^4 model we will meet the projection operators P_N , which do not commute with the paraproduct. Here we use a random operator technique from [GP15] to handle this operators. However, for the Φ_3^4 model this technique is not enough for our case and we have to estimate an additional error term D_N by stochastic calculations in Section 6.4 (see Remark 4.5).

Framework and main result

For $N \geq 1$, let $\Lambda_N = \{-N, -(N-1), \dots, N\}^3$. Set $\varepsilon = \frac{2}{2N+1}$. Every point $k \in \Lambda_N$ can be identified with $x = \varepsilon k \in \Lambda_\varepsilon = \{x = (x_1, x_2, x_3) \in \varepsilon \mathbb{Z}^3 : -1 < x_1, x_2, x_3 < 1\}$. We view Λ_ε as a discretisation of the continuous three-dimensional torus \mathbb{T}^3 identified with $[-1, 1]^3$. In the following for simplicity we fix a cylindrical Wiener process in (1.2) on $L^2(\mathbb{T}^3)$ given by $2^{-3/2} \sum_k \beta_k e^{i\pi k \cdot x}$ for $x \in \mathbb{T}^3$ and restrict it on $L^2(\Lambda_\varepsilon)$ as $W_N = 2^{-3/2} \sum_{|k|_\infty \leq N} \beta_k e^{i\pi k \cdot x}$ for $x \in \Lambda_\varepsilon$, which is also a cylindrical Wiener process on $L^2(\Lambda_\varepsilon)$. Here $\{\beta_k\}$ is a family of independent Brownian motions. Then for (1.4) and fixed N it is a finite dimensional SDE and we could easily obtain existence and uniqueness of solutions to (1.4) by [PR07], which implies that the solution to (1.4) has the same distribution as the solution to the following equation:

$$\begin{aligned} d\Phi^\varepsilon(t, x) &= (\Delta_\varepsilon \Phi^\varepsilon(t, x) - (\Phi^\varepsilon)^3(t, x) + (3C_0^\varepsilon - 9C_1^\varepsilon)\Phi^\varepsilon(t, x))dt \\ &\quad + dW_N(t, x) \\ \Phi^\varepsilon(0) &= \Phi_0^\varepsilon. \end{aligned} \tag{1.5}$$

Following [MW14] we discuss a suitable extension of functions defined on Λ_ε onto all of the torus \mathbb{T}^3 (which we identify with the interval $[-1, 1]^3$). For any function $Y : \Lambda_\varepsilon \rightarrow \mathbb{R}$, we define the discrete Fourier transform \hat{Y} through

$$\hat{Y}(k) = \begin{cases} \sum_{x \in \Lambda_\varepsilon} \varepsilon^3 Y(x) e^{-i\pi k \cdot x}, & \text{if } k \in \{-N, \dots, N\}^3 \\ 0 & \text{if } k \in \mathbb{Z}^3 \setminus \{-N, \dots, N\}^3. \end{cases}$$

In this context Fourier inversion states

$$Y(x) = \frac{1}{8} \sum_{k \in \mathbb{Z}^3} \hat{Y}(k) e^{i\pi k \cdot x} \text{ for all } x \in \Lambda_\varepsilon. \tag{1.6}$$

It is thus natural to extend Y to all of \mathbb{T}^3 by taking (1.6) as a definition of $Y(x)$ for $x \in \mathbb{T}^3 \setminus \Lambda_\varepsilon$. More explicitly, for $Y : \Lambda_\varepsilon \rightarrow \mathbb{R}$ we define $(\text{Ext}Y) : \mathbb{T}^3 \rightarrow \mathbb{R}$ as

$$\text{Ext}Y(x) = \frac{1}{2^3} \sum_{k \in \{-N, \dots, N\}^3} \sum_{y \in \Lambda_\varepsilon} \varepsilon^3 e^{i\pi k \cdot (x-y)} Y(y).$$

Let $P_t^\varepsilon = \text{Ext}e^{t\Delta_\varepsilon}$. By the definition of the operators Δ_ε we have

$$\widehat{e^{t\Delta_\varepsilon} v}(k) = \begin{cases} e^{-|k|^2 f(\varepsilon k)} \hat{v}(k), & \text{if } k \in \{-N, \dots, N\}^3 \\ 0 & \text{if } k \in \mathbb{Z}^3 \setminus \{-N, \dots, N\}^3. \end{cases}$$

Here

$$f(x) = \frac{4}{|x|^2} \left(\sin^2 \frac{x_1 \pi}{2} + \sin^2 \frac{x_2 \pi}{2} + \sin^2 \frac{x_3 \pi}{2} \right).$$

Now we extend the solution to all of \mathbb{T}^3 . In the following the Fourier transform and the inverse Fourier transform are denoted by \mathcal{F} and \mathcal{F}^{-1} . It is easy to see that

$$\text{Ext}\Phi^\varepsilon(t) = P_t^\varepsilon \text{Ext}\Phi_0^\varepsilon - \int_0^t P_{t-s}^\varepsilon Q_N [(\text{Ext}\Phi^\varepsilon)^3 - (3C_0^\varepsilon - 9C_1^\varepsilon)\text{Ext}\Phi^\varepsilon] ds + \int_0^t P_{t-s}^\varepsilon \text{Ext} dW_N, \tag{1.7}$$

where $Q_N u(x) = P_N u(x) + \Pi_N u(x)$ with

$$P_N = \mathcal{F}^{-1} 1_{|k|_\infty \leq N} \mathcal{F},$$

and Π_N is defined for u satisfying $\text{supp } \mathcal{F}u \subset \{k : |k|_\infty \leq 3N\}$

$$\begin{aligned} \Pi_N u(x) &= \sum_{i_1, i_2, i_3 \in \{-1, 0, 1\}, \sum_{j=1}^3 i_j^2 \neq 0} e_N^{i_1 i_2 i_3} \mathcal{F}^{-1} 1_{k \in P^{i_1 i_2 i_3}} \mathcal{F}u(x) \\ &= \sum_{i_1, i_2, i_3 \in \{-1, 0, 1\}, \sum_{j=1}^3 i_j^2 \neq 0} P_N [e_N^{i_1 i_2 i_3} u] \end{aligned}$$

with $P^{i_1 i_2 i_3} = \{k : k^j i_j > N \text{ if } i_j = -1, 1; |k^j| \leq N, \text{ if } i_j = 0\}$ is a rectangular division of $\mathbb{Z}^3 \setminus \{k \in \mathbb{Z}^3, |k|_\infty \leq N\}$, $e_N^{i_1 i_2 i_3} = \prod_{j=1}^3 e^{-i\pi(2N+1)i_j x^j}$ and $|k|_\infty = \max(|k^1|, |k^2|, |k^3|)$.

Now choose C_0^ε as in (6.3) and

$$C_1^\varepsilon = C_{11}^\varepsilon + \sum_{i_1, i_2, i_3 \in \{-1, 0, 1\}, \sum_{j=1}^3 i_j^2 \neq 0} C_{12}^{\varepsilon, i_1 i_2 i_3},$$

with $C_{11}^\varepsilon, C_{12}^{\varepsilon, i_1 i_2 i_3}$ as in (6.4) and (6.5) respectively. In the following we omit the summation with respect to i_1, i_2, i_3 if there's no confusion.

The main result of this paper is the following theorem:

Theorem 1.1 Let $z \in (1/2, 2/3)$ and $\Phi_0 \in \mathcal{C}^{-z}$. Let (Φ, τ) be the unique (maximal in time) solution to (1.2) and let for $\varepsilon \in (0, 1)$ the function Φ^ε be the unique solution to (1.5) on $[0, \infty)$. If the initial data satisfy $\text{Ext}\Phi_0^\varepsilon - \Phi_0 \rightarrow 0$ in \mathcal{C}^{-z} then there exists a sequence of random time τ_L such that $\lim_{L \rightarrow \infty} \tau_L = \tau$ and

$$\sup_{t \in [0, \tau_L]} \|\text{Ext}\Phi^\varepsilon - \Phi\|_{-z} \rightarrow 0 \quad \text{in probability, as } \varepsilon \rightarrow 0.$$

Remark 1.2 (1) Existence and uniqueness of (Φ, τ) has been obtained in [Hai14, CC13]. For the definition of \mathcal{C}^{-z} and norm $\|\cdot\|_{-z}$ see Section 2.

(2) The constant C_1^ε is divided into two parts: C_{11}^ε and C_{12}^ε which correspond to terms with P_N and Π_N respectively. In fact $C_0^\varepsilon \simeq \frac{1}{\varepsilon}$, $C_{11}^\varepsilon \simeq -\log \varepsilon$ and $C_{12}^{\varepsilon, i_1 i_2 i_3} \simeq 1$.

The structure of the paper is organized as follows. In Section 2, we recall some basic notions and results for the paracontrolled distribution method. In Section 3 we prove some estimates for the approximating operators. In Section 4 we use the paracontrolled distribution method to prove uniform bounds for the lattice approximation equations. In Section 5 we give the proof of our main result. In Section 6 convergence of several stochastic terms is proved.

2 Besov spaces and paraproduct

In the following we recall the definitions and some properties of Besov spaces and paraproducts. For a general introduction to these theories we refer to [BCD11, GIP13]. First we introduce

the following notations. The space of real valued infinitely differentiable functions of compact support is denoted by $\mathcal{D}(\mathbb{R}^d)$ or \mathcal{D} . The space of Schwartz functions is denoted by $\mathcal{S}(\mathbb{R}^d)$. Its dual, the space of tempered distributions is denoted by $\mathcal{S}'(\mathbb{R}^d)$.

Let $\chi, \theta \in \mathcal{D}$ be nonnegative radial functions on \mathbb{R}^d , such that

- i. the support of χ is contained in a ball and the support of θ is contained in an annulus;
- ii. $\chi(z) + \sum_{j \geq 0} \theta(2^{-j}z) = 1$ for all $z \in \mathbb{R}^d$.
- iii. $\text{supp}(\chi) \cap \text{supp}(\theta(2^{-j}\cdot)) = \emptyset$ for $j \geq 1$ and $\text{supp}(\theta(2^{-i}\cdot)) \cap \text{supp}(\theta(2^{-j}\cdot)) = \emptyset$ for $|i-j| > 1$.

We call such (χ, θ) dyadic partition of unity, and for the existence of dyadic partitions of unity we refer to [BCD11, Proposition 2.10]. The Littlewood-Paley blocks are now defined as

$$\Delta_{-1}u = \mathcal{F}^{-1}(\chi \mathcal{F}u) \quad \Delta_j u = \mathcal{F}^{-1}(\theta(2^{-j}\cdot) \mathcal{F}u).$$

For $\alpha \in \mathbb{R}$, the Hölder-Besov space \mathcal{C}^α is given by $\mathcal{C}^\alpha = B_{\infty, \infty}^\alpha(\mathbb{R}^d, \mathbb{R}^n)$, where for $p, q \in [1, \infty]$ we define

$$B_{p,q}^\alpha(\mathbb{R}^d) = \{u \in \mathcal{S}'(\mathbb{R}^d) : \|u\|_{B_{p,q}^\alpha} = (\sum_{j \geq -1} (2^{j\alpha} \|\Delta_j u\|_{L^p})^q)^{1/q} < \infty\},$$

with the usual interpretation as l^∞ norm in case $q = \infty$. We write $\|\cdot\|_\alpha$ instead of $\|\cdot\|_{B_{\infty, \infty}^\alpha}$ in the following for simplicity. We also use $C_T E$ to denote $C([0, T]; E)$.

We point out that everything above and everything that follows can be applied to distributions on the torus (see [S85, SW71]). More precisely, let $\mathcal{S}'(\mathbb{T}^d)$ be the space of distributions on \mathbb{T}^d . Therefore, Besov spaces on the torus with general indices $p, q \in [1, \infty]$ are defined as

$$B_{p,q}^\alpha(\mathbb{T}^d) = \{u \in \mathcal{S}'(\mathbb{T}^d) : \|u\|_{B_{p,q}^\alpha} = (\sum_{j \geq -1} (2^{j\alpha} \|\Delta_j u\|_{L^p(\mathbb{T}^d)})^q)^{1/q} < \infty\}.$$

We will need the following Besov embedding theorem on the torus (c.f. [GIP13, Lemma 41]):

Lemma 2.1 Let $1 \leq p_1 \leq p_2 \leq \infty$ and $1 \leq q_1 \leq q_2 \leq \infty$, and let $\alpha \in \mathbb{R}$. Then $B_{p_1, q_1}^\alpha(\mathbb{T}^d)$ is continuously embedded in $B_{p_2, q_2}^{\alpha - d(1/p_1 - 1/p_2)}(\mathbb{T}^d)$.

Now we recall the following paraproduct introduced by Bony (see [Bon81]). In general, the product fg of two distributions $f \in \mathcal{C}^\alpha, g \in \mathcal{C}^\beta$ is well defined if and only if $\alpha + \beta > 0$. In terms of Littlewood-Paley blocks, the product fg can be formally decomposed as

$$fg = \sum_{j \geq -1} \sum_{i \geq -1} \Delta_i f \Delta_j g = \pi_{<}(f, g) + \pi_0(f, g) + \pi_{>}(f, g),$$

with

$$\pi_{<}(f, g) = \pi_{>}(g, f) = \sum_{j \geq -1} \sum_{i < j-1} \Delta_i f \Delta_j g, \quad \pi_0(f, g) = \sum_{|i-j| \leq 1} \Delta_i f \Delta_j g.$$

We also use the notation for $j \geq 0$

$$S_j f = \sum_{i \leq j-1} \Delta_i f.$$

Moreover define

$$\psi_{<}(k_1, k_2) = \sum_{j \geq -1} \sum_{i < j-1} \theta(2^{-i}k_1) \theta(2^{-j}k_2)$$

and

$$\psi_0(k_1, k_2) = \sum_{|i-j| \leq 1} \theta(2^{-i}k_1)\theta(2^{-j}k_2).$$

We will use without comment that $\|\cdot\|_\alpha \leq \|\cdot\|_\beta$ for $\alpha \leq \beta$, that $\|\cdot\|_{L^\infty} \lesssim \|\cdot\|_\alpha$ for $\alpha > 0$, and that $\|\cdot\|_\alpha \lesssim \|\cdot\|_{L^\infty}$ for $\alpha \leq 0$. We will also use that $\|S_j u\|_{L^\infty} \lesssim 2^{-j\alpha} \|u\|_\alpha$ for $\alpha < 0, j \geq 0$ and $u \in \mathcal{C}^\alpha$, where $\|\cdot\|_\alpha$ denotes the norm in $\mathcal{C}^\alpha, \alpha \in \mathbb{R}$.

The basic result about these bilinear operations is given by the following estimates:

Lemma 2.2 (Paraproduct estimates, [Bon 81, GIP13, Lemma 2]) For any $\beta \in \mathbb{R}$ we have

$$\|\pi_{<}(f, g)\|_\beta \lesssim \|f\|_{L^\infty} \|g\|_\beta \quad f \in L^\infty, g \in \mathcal{C}^\beta,$$

and for $\alpha < 0$ furthermore

$$\|\pi_{<}(f, g)\|_{\alpha+\beta} \lesssim \|f\|_\alpha \|g\|_\beta \quad f \in \mathcal{C}^\alpha, g \in \mathcal{C}^\beta.$$

For $\alpha + \beta > 0$ we have

$$\|\pi_0(f, g)\|_{\alpha+\beta} \lesssim \|f\|_\alpha \|g\|_\beta \quad f \in \mathcal{C}^\alpha, g \in \mathcal{C}^\beta.$$

The following basic commutator lemma is important for our use:

Lemma 2.3 ([GIP13, Lemma 5]) Assume that $\alpha \in (0, 1)$ and $\beta, \gamma \in \mathbb{R}$ are such that $\alpha + \beta + \gamma > 0$ and $\beta + \gamma < 0$. Then for smooth f, g, h , the trilinear operator

$$C(f, g, h) = \pi_0(\pi_{<}(f, g), h) - f\pi_0(g, h)$$

allows for the bound

$$\|C(f, g, h)\|_{\alpha+\beta+\gamma} \lesssim \|f\|_\alpha \|g\|_\beta \|h\|_\gamma.$$

Thus, C can be uniquely extended to a bounded trilinear operator from $\mathcal{C}^\alpha \times \mathcal{C}^\beta \times \mathcal{C}^\gamma$ to $\mathcal{C}^{\alpha+\beta+\gamma}$.

Now we recall the following estimate for heat semigroup $P_t := e^{t\Delta}$.

Lemma 2.4 ([GIP13, Lemma 47]) Let $u \in \mathcal{C}^\alpha$ for some $\alpha \in \mathbb{R}$. Then for every $\delta \geq 0$

$$\|P_t u\|_{\alpha+\delta} \lesssim t^{-\delta/2} \|u\|_\alpha.$$

Lemma 2.5 ([CC13, Lemma A.1]) Let $u \in \mathcal{C}^\alpha$ for some $\alpha < 1$ and $v \in \mathcal{C}^\beta$ for some $\beta \in \mathbb{R}$. Then for $\delta \geq \alpha + \beta$

$$\|P_t \pi_{<}(u, v) - \pi_{<}(u, P_t v)\|_\delta \lesssim t^{\frac{\alpha+\beta-\delta}{2}} \|u\|_\alpha \|v\|_\beta.$$

Lemma 2.6 ([CC13, Lemma 2.5]) Let $u \in \mathcal{C}^{\alpha+\delta}$ for some $\alpha \in \mathbb{R}, \delta > 0$. Then for every $t \geq 0$

$$\|(P_t - I)u\|_\alpha \lesssim t^{\delta/2} \|u\|_{\alpha+\delta}.$$

We also have the following result.

Lemma 2.7 (Bernstein type lemma) Let $u \in \mathcal{C}^\alpha$ for some $\alpha \in \mathbb{R}$.

1) If $\text{supp}\mathcal{F}u \subset \{k : |k| \leq CN\}$ for some $C > 0$ then for $\beta > \alpha$

$$\|u\|_\beta \lesssim N^{\beta-\alpha} \|u\|_\alpha.$$

2) If $\text{supp}\mathcal{F}u \subset \{k : |k| > CN\}$ for some $C > 0$ then for $\beta < \alpha$

$$\|u\|_\beta \lesssim N^{\beta-\alpha} \|u\|_\alpha.$$

Here all the constants we omit are independent of N .

Proof We have

$$\|u\|_\beta = \sup_j 2^{j\beta} \|\Delta_j u\|_{L^\infty} = \sup_j 2^{j(\beta-\alpha)} 2^{j\alpha} \|\Delta_j u\|_{L^\infty}.$$

For the first case we have $\Delta_j u \neq 0$ iff $2^j \lesssim N$, which implies the first result. If $\text{supp}\mathcal{F}u \subset \{k : |k| > CN\}$ we have $\Delta_j u \neq 0$ iff $2^j \gtrsim N$ which implies the second result. \square

3 Estimates for the approximated operators

Now we prove the following estimates for the approximated operators on \mathbb{T}^3 . First we consider estimate for P_N and Π_N :

Lemma 3.1 Let $u \in \mathcal{C}^\alpha$ for some $\alpha \in \mathbb{R}$. Then for any $\kappa > 0$ small enough we have the following estimate:

(1) (Estimate for P_N)

$$\|P_N u\|_{\alpha-\kappa} \lesssim \|u\|_\alpha, \quad \|(I - P_N)u\|_{\alpha-\kappa} \lesssim N^{-\kappa/2} \|u\|_\alpha.$$

(2) (Estimate for Π_N) If $\alpha > 0$ then for u satisfying $\text{supp}\mathcal{F}u \subset \{k : |k|_\infty \leq 3N\}$

$$\|\Pi_N u\|_{\alpha-\kappa} \lesssim N^{-\kappa/2} \|u\|_\alpha.$$

If $\alpha < 0$ and $\text{supp}\mathcal{F}u \subset \{k : |k|_\infty \leq N\}$ then

$$\|e_N^{i_1 i_2 i_3} u\|_{\alpha-\kappa} \lesssim N^{-\kappa/2} \|u\|_\alpha.$$

Here all the constants we omit are independent of N .

Proof We have for p large enough

$$\|P_N u\|_{\alpha-\kappa} \lesssim \|P_N u\|_{B_{p,\infty}^\alpha} \lesssim \|u\|_{B_{p,\infty}^\alpha} \lesssim \|u\|_\alpha,$$

where in the first inequality we used Lemma 2.1 and in the second inequality we used that $1_{|k|_\infty \leq N}$ is an L^p multiplier. Similarly

$$\|(I - P_N)u\|_{\alpha-\kappa} \lesssim N^{-\kappa/2} \|(I - P_N)u\|_{\alpha-\kappa/2} \lesssim N^{-\kappa/2} \|u\|_\alpha,$$

where in the first inequality we used Lemma 2.7 and in the second inequality we used the result for P_N . Moreover for $\alpha > 5\kappa/4$

$$\begin{aligned} \|\Pi_N u\|_{\alpha-\kappa} &\lesssim N^{\alpha-5\kappa/4} \|\Pi_N u\|_{\kappa/4} \lesssim N^{\alpha-\kappa} \|\mathcal{F}^{-1} 1_{k \in P^{i_1 i_2 i_3}} \mathcal{F} u\|_{\kappa/4} \\ &\lesssim N^{-\kappa/2} \|\mathcal{F}^{-1} 1_{k \in P^{i_1 i_2 i_3}} \mathcal{F} u\|_{\alpha-\kappa/4} \lesssim N^{-\kappa/2} \|u\|_\alpha. \end{aligned}$$

Here in the first and third inequalities we used Lemma 2.7, in the second inequality we used that $\|e_N^{i_1 i_2 i_3}\|_{\kappa/2} \lesssim N^{\kappa/2}$ and in the last inequality we used similar argument for P_N since $1_{k \in P^{i_1 i_2 i_3}}$ is an L^p multiplier. Similarly for $\alpha < 0$

$$\|e_N^{i_1 i_2 i_3} u\|_{\alpha-\kappa} \lesssim N^{\alpha-3\kappa/2} \|e_N^{i_1 i_2 i_3} u\|_{\kappa/2} \lesssim N^{\alpha-\kappa} \|u\|_{\kappa/2} \lesssim N^{-\kappa/2} \|u\|_{\alpha}.$$

Here in the first inequality we used $\text{supp} \mathcal{F}(e_N^{i_1 i_2 i_3} u) \subset \{k : |k| > N\}$ and Lemma 2.7. Thus the result follows. \square

Now want to prove estimates for $P_t^\varepsilon = P_N e^{t\Delta_\varepsilon}$. In fact,

$$P_t^\varepsilon = \mathcal{F}^{-1} 1_{|k|_\infty \leq N} e^{-t|k|^2 f(\varepsilon k)} \varphi(\varepsilon k) \mathcal{F} = P_N \tilde{P}_t^\varepsilon,$$

with

$$\tilde{P}_t^\varepsilon := \mathcal{F}^{-1} e^{-t|k|^2 f(\varepsilon k)} \varphi(\varepsilon k) \mathcal{F},$$

where φ is a smooth function and equals 1 on $\{|x|_\infty \leq 1\}$ with support in $\{|x| \leq 1.8\}$. Then by a similar argument as [GIP13, Lemma 47] we have the following result:

Lemma 3.2 Let $u \in \mathcal{C}^\alpha$ for some $\alpha \in \mathbb{R}$. Then for every $\delta \geq 0, \kappa > 0, t > 0$,

$$\|P_t^\varepsilon u\|_{\alpha+\delta-\kappa} \lesssim t^{-\delta/2} \|u\|_{\alpha},$$

$$\|(P_t^\varepsilon - P_t)u\|_{\alpha+\delta-\kappa} \lesssim \varepsilon^{\kappa/2} t^{-\delta/2} \|u\|_{\alpha}.$$

Here the constants we omit are independent of N .

Proof For the first result by Lemma 3.1 it suffices to prove

$$\|\tilde{P}_t^\varepsilon u\|_{\alpha+\delta} \lesssim t^{-\delta/2} \|u\|_{\alpha}. \quad (3.1)$$

In the following we consider (3.1) and have for $j \geq 0$

$$\begin{aligned} \|\Delta_j \tilde{P}_t^\varepsilon u\|_{L^\infty} &= \|\mathcal{F}^{-1} \theta_j \phi^\varepsilon \mathcal{F} u\|_{L^\infty} = \|\mathcal{F}^{-1} \theta_j \tilde{\theta}(2^{-j} \cdot) \phi^\varepsilon \mathcal{F} u\|_{L^\infty} \\ &\leq \|\mathcal{F}^{-1}(\phi^\varepsilon \tilde{\theta}(2^{-j} \cdot))\|_{L^1} \|\Delta_j u\|_{L^\infty}. \end{aligned}$$

Here

$$\phi^\varepsilon(\xi) = e^{-t|\xi|^2 f(\varepsilon \xi)} \varphi(\varepsilon \xi),$$

and $\tilde{\theta}$ be a smooth function supported in an annulus such that $\tilde{\theta}\theta = \theta$. Then we get that for $\delta \geq 0$

$$\begin{aligned} \|\mathcal{F}^{-1}(\phi^\varepsilon \tilde{\theta}(2^{-j} \cdot))\|_{L^1} &= \|\mathcal{F}^{-1}(\phi^\varepsilon(2^j \cdot) \tilde{\theta})\|_{L^1} \lesssim \|(1 - \Delta)^2(\phi^\varepsilon(2^j \cdot) \tilde{\theta})\|_{L^1} \\ &\lesssim \sum_{0 \leq |k| \leq 4} 2^{j|k|} \|(D_k \phi^\varepsilon)(2^j \cdot)\|_{L^\infty} \lesssim \sum_{0 \leq |k| \leq 4} 2^{j|k|} \frac{1}{2^{j|k|} (2^j \sqrt{t})^\delta} \lesssim (2^j \sqrt{t})^{-\delta}. \end{aligned}$$

Here in the third inequality we used that $f(\varepsilon \xi) \geq c > 0$ and $|\varepsilon \xi| \lesssim 1$ on the support of φ^ε which implies that for any multiindices k satisfying $|k| \leq 4$ and every $\delta \geq 0$ $|D_k \phi^\varepsilon(\xi)| \lesssim \frac{1}{|\xi|^{|k|+\delta} t^{\delta/2}}$.

For $j = -1$ we can use Bernstein's lemma to obtain the estimate. Thus (3.1) follows.

For the second result we have

$$P_t^\varepsilon - P_t = P_N(\tilde{P}_t^\varepsilon - P_t) + (I - P_N)P_t.$$

By Lemmas 2.4 and 3.1 it is sufficient to consider $\tilde{P}_t^\varepsilon - P_t$. Since $\phi^\varepsilon(\xi) - \phi(\xi) = \varphi(\varepsilon\xi)(e^{-t|\xi|^2 f(\varepsilon\xi)} - e^{-t|\xi|^2}) + (\varphi(\varepsilon\xi) - 1)e^{-t|\xi|^2}$ and $|\varphi(\varepsilon\xi) - 1| \lesssim |\varepsilon\xi|^\delta$, $|f(\varepsilon\xi) - 1| \lesssim |\varepsilon\xi|^\delta$, we obtain that for any multiindices k satisfying $|k| \leq 4$ and every $\delta \geq 0$, $0 < \eta < 1$ $|D_k(\phi^\varepsilon - \phi)(\xi)| \leq \frac{(\varepsilon|\xi|)^\eta}{|\xi|^{|k|+\delta t^{\delta/2}}}$. Thus the second result follows by a similar argument as the calculation for (3.1). \square

Now we prove a commutator estimate for P_t^ε . However P_N does not commute with para-product, we could only obtain the commutator estimate for \tilde{P}_t^ε .

Lemma 3.3 Let $u \in \mathcal{C}^\alpha$ for some $\alpha < 1$ and $v \in \mathcal{C}^\beta$ for some $\beta \in \mathbb{R}$. Then for $\delta \geq \alpha + \beta$ and any $\kappa > 0$

$$\|P_t^\varepsilon \pi_{<}(u, v) - P_N \pi_{<}(u, \tilde{P}_t^\varepsilon v)\|_{\delta-\kappa} \lesssim t^{\frac{\alpha+\beta-\delta}{2}} \|u\|_\alpha \|v\|_\beta, \quad (3.2)$$

$$\|(P_t^\varepsilon - P_t) \pi_{<}(u, v) - P_N \pi_{<}(u, \tilde{P}_t^\varepsilon v) - \pi_{<}(u, P_t v)\|_{\delta-\kappa} \lesssim \varepsilon^{\kappa/2} t^{\frac{\alpha+\beta-\delta}{2}} \|u\|_\alpha \|v\|_\beta. \quad (3.3)$$

Here the constants we omit are independent of N .

Proof We have

$$P_t^\varepsilon \pi_{<}(u, v) - P_N \pi_{<}(u, \tilde{P}_t^\varepsilon v) = P_N (\tilde{P}_t^\varepsilon \pi_{<}(u, v) - \pi_{<}(u, \tilde{P}_t^\varepsilon v)).$$

By Lemma 3.1 it suffices to prove that

$$\|\tilde{P}_t^\varepsilon \pi_{<}(u, v) - \pi_{<}(u, \tilde{P}_t^\varepsilon v)\|_\delta \lesssim t^{\frac{\alpha+\beta-\delta}{2}} \|u\|_\alpha \|v\|_\beta. \quad (3.4)$$

In fact, we have

$$\tilde{P}_t^\varepsilon \pi_{<}(u, v) - \pi_{<}(u, \tilde{P}_t^\varepsilon v) = \sum_{j=-1}^{\infty} (\tilde{P}_t^\varepsilon (S_{j-1} u \Delta_j v) - S_{j-1} u \tilde{P}_t^\varepsilon \Delta_j v).$$

We also have that the Fourier transform of $\tilde{P}_t^\varepsilon (S_{j-1} u \Delta_j v) - S_{j-1} u \tilde{P}_t^\varepsilon \Delta_j v$ has its support in a suitable annulus $2^j \mathcal{A}$. Let $\psi \in \mathcal{D}(\mathbb{R}^3)$ with support in an annulus $\tilde{\mathcal{A}}$ be such that $\psi = 1$ on \mathcal{A} .

Thus by the same argument as the proof of [CC13, Lemma A.1] we obtain that

$$\begin{aligned} & \|[(\psi(2^{-j}\cdot)\phi^\varepsilon)(D), S_{j-1}u]\Delta_j v\|_{L^\infty} \\ & \lesssim \sum_{\eta \in \mathbb{N}^d, |\eta|=1} \|x^\eta \mathcal{F}^{-1}(\psi(2^{-j}\cdot)\phi^\varepsilon)\|_{L^1} \|\partial^\eta S_{j-1}u\|_{L^\infty} \|\Delta_j v\|_{L^\infty}. \end{aligned}$$

Here $(\psi(2^{-j}\cdot)\phi^\varepsilon)(D)u = \mathcal{F}^{-1}(\psi(2^{-j}\cdot)\phi^\varepsilon \mathcal{F}u)$, $[(\psi(2^{-j}\cdot)\phi^\varepsilon)(D), S_{j-1}u]$ denotes the commutator.

Now we have that

$$\begin{aligned} & \|x^\eta \mathcal{F}^{-1}(\psi(2^{-j}\cdot)\phi^\varepsilon)\|_{L^1} \\ & \leq 2^{-j} \|\mathcal{F}^{-1}(\partial^\eta \psi)(2^{-j}\cdot)\phi^\varepsilon\|_{L^1} + \|\mathcal{F}^{-1}(\psi(2^{-j}\cdot)\partial^\eta \phi^\varepsilon)\|_{L^1} \\ & = 2^{-j} \|\mathcal{F}^{-1}(\partial^\eta \psi(\cdot)\phi^\varepsilon(2^j\cdot))\|_{L^1} + \|\mathcal{F}^{-1}(\psi(\cdot)\partial^\eta \phi^\varepsilon(2^j\cdot))\|_{L^1} \\ & \lesssim 2^{-j} \|(1 + |\cdot|^2)^2 \mathcal{F}^{-1}(\partial^\eta \psi(\cdot)\phi^\varepsilon(2^j\cdot))\|_{L^\infty} + \|(1 + |\cdot|^2)^2 \mathcal{F}^{-1}(\psi(\cdot)\partial^\eta \phi^\varepsilon(2^j\cdot))\|_{L^\infty} \\ & = 2^{-j} \|\mathcal{F}^{-1}((1 - \Delta)^2(\partial^\eta \psi(\cdot)\phi^\varepsilon(2^j\cdot)))\|_{L^\infty} + \|\mathcal{F}^{-1}((1 - \Delta)^2(\psi(\cdot)\partial^\eta \phi^\varepsilon(2^j\cdot)))\|_{L^\infty} \\ & \lesssim 2^{-j} \|(1 - \Delta)^2(\partial^\eta \psi(\cdot)\phi^\varepsilon(2^j\cdot))\|_{L^1} + \|(1 - \Delta)^2(\psi(\cdot)\partial^\eta \phi^\varepsilon(2^j\cdot))\|_{L^1} \\ & \lesssim 2^{-j} \sum_{0 \leq |m| \leq 4} (2^j)^{|m|} \frac{t^{-\mu} 2^{-2j\mu}}{(2^j)^{|m|}} + \sum_{|m| \leq 5} (2^j)^{|m|} \frac{t^{-\mu} 2^{-2j\mu}}{(2^j)^{|m|+1}} \\ & \lesssim 2^{-j} t^{-\mu} 2^{-2j\mu}, \end{aligned}$$

where in the fourth inequality we used $|D^m \phi^\varepsilon(\xi)| \lesssim |\xi|^{-|m|} t^{-\mu} |\xi|^{-2\mu}$, $\mu \geq 0$ for any multiindices m satisfying $|m| \leq 5$. Hence we get that

$$\|[\psi(2^{-j} \cdot) \phi^\varepsilon(D), S_{j-1} u] \Delta_j v\|_{L^\infty} \lesssim t^{\frac{\alpha+\beta-\delta}{2}} 2^{j(\alpha+\beta-\delta)} 2^{-j(\alpha+\beta)} \|u\|_\alpha \|v\|_\beta,$$

which yields (3.4) by the same argument as in the proof of [CC13, Lemma A.1].

Moreover we have

$$\begin{aligned} & (P_t^\varepsilon - P_t) \pi_<(u, v) - P_N \pi_<(u, \tilde{P}_t^\varepsilon v) - \pi_<(u, P_t v) \\ &= P_N [(\tilde{P}_t^\varepsilon - P_t) \pi_<(u, v) - \pi_<(u, (\tilde{P}_t^\varepsilon - P_t) v)] - (I - P_N) (P_t \pi_<(u, v) - \pi_<(u, P_t v)). \end{aligned}$$

The estimate for the second term can be obtained by Lemmas 2.5, 2.7 and 3.1. By a similar argument as Lemma 3.2 we obtain that for any multiindices k satisfying $|k| \leq 5$ and every $\delta \geq 0, 0 < \eta < 1$ $|D_k(\phi^\varepsilon - \phi)(\xi)| \leq \frac{(\varepsilon|\xi|)^\eta}{|\xi|^{|k|+\delta t^{\delta/2}}}$. Thus (3.3) follows by a similar argument as the proof of (3.4). \square

Now we prove the following continuity result for P_t^ε .

Lemma 3.4 Let $u \in \mathcal{C}^{\alpha+\delta}$ for some $\alpha \in \mathbb{R}, 0 < \delta < 1$. Then for every $\varepsilon \in (0, 1), \kappa > 0, t > s > 0$

$$\|(P_t^\varepsilon - P_s^\varepsilon)u\|_{\alpha-\kappa} \lesssim (t-s)^{\delta/2} \|u\|_{\alpha+\delta}.$$

Here the constants are independent of N .

Proof We have $(P_t^\varepsilon - P_s^\varepsilon)u = P_N(\tilde{P}_t^\varepsilon - \tilde{P}_s^\varepsilon)u$. By Lemma 3.1 it suffices to prove that

$$\|(\tilde{P}_t^\varepsilon - \tilde{P}_s^\varepsilon)u\|_\alpha \lesssim (t-s)^{\delta/2} \|u\|_{\alpha+\delta}.$$

By $|1 - e^{-(t-s)f(\varepsilon\xi)|\xi|^2}| \leq (t-s)^{\delta/2} |\xi|^\delta$ we obtain that for any multiindices k satisfying $|k| \leq 4$ and any $\delta \geq 0$ $|D_k(\phi_t^\varepsilon - \phi_s^\varepsilon)(\xi)| \lesssim \frac{(t-s)^{\delta/2} |\xi|^\delta}{|\xi|^{|k|}}$. Thus by a similar argument as Lemma 3.2 the result follows. \square

4 Paracontrolled analysis for the approximating equations

Now let $u^\varepsilon = \text{Ext}\Phi^\varepsilon$ for simplicity and we have the following equation:

$$u^\varepsilon(t) = P_t^\varepsilon \text{Ext}\Phi_0^\varepsilon - \int_0^t P_{t-s}^\varepsilon Q_N[(u^\varepsilon)^3 - (3C_0^\varepsilon - 9C_1^\varepsilon)u^\varepsilon] ds + \int_0^t P_{t-s}^\varepsilon P_N dW. \quad (4.1)$$

Therefore it suffices to prove the convergence result for solutions to (4.1). In this section we give an uniform estimate for solutions to (4.1) by using paracontrolled analysis.

In this section we fix $\delta, \beta, \kappa, \gamma > 0$ satisfying

$$2z - 1 \geq \delta > 2\kappa, \quad \beta > \frac{\delta}{2}, \quad \beta + \frac{\delta}{2} + \kappa < \gamma, \quad 5\kappa + \frac{\delta}{2} + \beta + 3\gamma < 2 - 3z.$$

Now we split (4.1) into the following three equations:

$$u_1^\varepsilon = \int_{-\infty}^t P_{t-s}^\varepsilon P_N dW,$$

$$u_2^\varepsilon = - \int_0^t P_{t-s}^\varepsilon Q_N[(u_1^\varepsilon)^{\diamond,3}] ds$$

and

$$u_3^\varepsilon(t) = P_t^\varepsilon(\text{Ext}\Phi_0^\varepsilon - u_1^\varepsilon(0)) - \int_0^t P_{t-s}^\varepsilon \left[Q_N[6u_1^\varepsilon \diamond u_2^\varepsilon u_3^\varepsilon + 3u_1^\varepsilon (u_3^\varepsilon)^2 + 3u_1^\varepsilon \diamond (u_2^\varepsilon)^2 + (u_2^\varepsilon + u_3^\varepsilon)^3] \right. \\ \left. + P_N[3(u_1^\varepsilon)^{\diamond,2} \diamond (u_2^\varepsilon + u_3^\varepsilon) + 3e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2} \diamond (u_2^\varepsilon + u_3^\varepsilon) - 9\varphi^\varepsilon u^\varepsilon] \right] ds. \quad (4.2)$$

Here

$$\begin{aligned} (u_1^\varepsilon)^{\diamond,2} &:= (u_1^\varepsilon)^2 - C_0^\varepsilon, \\ (u_1^\varepsilon)^{\diamond,3} &:= (u_1^\varepsilon)^3 - 3C_0^\varepsilon u_1^\varepsilon, \\ u_1^\varepsilon \diamond u_2^\varepsilon &:= u_2^\varepsilon u_1^\varepsilon, \\ u_1^\varepsilon \diamond (u_2^\varepsilon)^2 &:= (u_2^\varepsilon)^2 u_1^\varepsilon, \\ (u_1^\varepsilon)^{\diamond,2} \diamond u_2^\varepsilon &:= \pi_{<}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) + \pi_{>}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) + \pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) \\ &= u_2^\varepsilon (u_1^\varepsilon)^{\diamond,2} + 3(C_{11}^\varepsilon + \varphi_1^\varepsilon) u_1^\varepsilon, \\ (u_1^\varepsilon)^{\diamond,2} \diamond u_3^\varepsilon &:= \pi_{<}(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) + \pi_{>}(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) + \pi_{0,\diamond}(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) \\ &= u_3^\varepsilon (u_1^\varepsilon)^{\diamond,2} + 3(C_{11}^\varepsilon + \varphi_1^\varepsilon)(u_2^\varepsilon + u_3^\varepsilon), \\ e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2} \diamond u_2^\varepsilon &:= \pi_{<}(u_2^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2}) + \pi_{>}(u_2^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2}) + \pi_{0,\diamond}(u_2^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2}) \\ &= e_N^{i_1 i_2 i_3} u_2^\varepsilon (u_1^\varepsilon)^{\diamond,2} + 3(C_{12}^{\varepsilon, i_1 i_2 i_3} + \varphi_2^{\varepsilon, i_1 i_2 i_3}) u_1^\varepsilon, \\ e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2} \diamond u_3^\varepsilon &:= \pi_{<}(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond,2} e_N^{i_1 i_2 i_3}) + \pi_{>}(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond,2} e_N^{i_1 i_2 i_3}) + \pi_{0,\diamond}(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond,2} e_N^{i_1 i_2 i_3}) \\ &= u_3^\varepsilon (u_1^\varepsilon)^{\diamond,2} e_N^{i_1 i_2 i_3} + 3(C_{12}^{\varepsilon, i_1 i_2 i_3} + \varphi_2^{\varepsilon, i_1 i_2 i_3})(u_2^\varepsilon + u_3^\varepsilon), \\ \varphi^\varepsilon &:= \varphi_1^\varepsilon + \varphi_2^\varepsilon, \end{aligned}$$

where $C_0^\varepsilon, C_{1i}^\varepsilon, \varphi_i^\varepsilon$ are defined in Section 6. Moreover there exist $\varphi \in C((0, T]; \mathbb{R})$ such that for $\rho > 0$ small enough $\sup_{t \in [0, T]} t^\rho |\varphi^\varepsilon - \varphi| \rightarrow 0$ as $\varepsilon \rightarrow 0$.

Define

$$K^\varepsilon(t) := \int_0^t P_{t-s}^\varepsilon (u_1^\varepsilon)^{\diamond,2} ds, \quad \tilde{K}^\varepsilon(t) := \int_0^t \tilde{P}_{t-s}^\varepsilon (u_1^\varepsilon)^{\diamond,2} ds,$$

and

$$K_1^\varepsilon(t) := \int_0^t P_{t-s}^\varepsilon [e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2}] ds, \quad \tilde{K}_1^\varepsilon(t) := \int_0^t \tilde{P}_{t-s}^\varepsilon [e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2}] ds.$$

Also define

$$\pi_{0,\diamond}(K^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) := \pi_0(K^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) - C_{11}^\varepsilon - \varphi_1^\varepsilon,$$

and

$$\pi_{0,\diamond}(K_1^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2}) := \pi_0(K_1^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2}) - C_{12}^{\varepsilon, i_1 i_2 i_3} - \varphi_2^{\varepsilon, i_1 i_2 i_3}.$$

Now we introduce the following notations:

$$C_W^\varepsilon(T) := \sup_{t \in [0, T]} (\|u_1^\varepsilon\|_{-1/2-\delta/2} + \|(u_1^\varepsilon)^{\diamond, 2}\|_{-1-\delta/2} + \|u_2^\varepsilon\|_{1/2-\delta} + \|\pi_0(u_2^\varepsilon, u_1^\varepsilon)\|_{-\delta} \\ + \|\pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond, 2})\|_{-1/2-\delta} + \|\pi_{0,\diamond}(K^\varepsilon, (u_1^\varepsilon)^{\diamond, 2})\|_{-\delta} + \|u_2^\varepsilon\|_{C_T^{1/4-\delta-\kappa/2} C^{\kappa/2}},$$

and

$$E_W^\varepsilon(T) := \sup_{t \in [0, T]} (\|(u_1^\varepsilon)^{\diamond, 2} e_N^{i_1 i_2 i_3}\|_{-1-\delta/2} + \|\pi_0(u_2^\varepsilon, e_N^{i_1 i_2 i_3} u_1^\varepsilon)\|_{-\delta} + \|\pi_{0,\diamond}(u_2^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})\|_{-1/2-\delta} \\ + \|\pi_0(K^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})\|_{-\delta} + \|\pi_0(K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2})\|_{-\delta} + \|\pi_{0,\diamond}(K_1^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})\|_{-\delta}).$$

Here E_W^ε appears as an error term for the lattice approximations which goes to 0 in probability (see Section 6.2).

Then Lemma 3.2 and (3.1) implies that for $t \in [0, T]$

$$\|K^\varepsilon(t)\|_{1-\delta} + \|\tilde{K}^\varepsilon(t)\|_{1-\delta} \lesssim t^{\delta/4} C_W^\varepsilon, \quad \|K_1^\varepsilon(t)\|_{1-\delta} + \|\tilde{K}_1^\varepsilon(t)\|_{1-\delta} \lesssim t^{\delta/4} E_W^\varepsilon. \quad (4.3)$$

Now we could write the paracontrolled ansatz as follows:

$$u_3^\varepsilon = -3P_N[\pi_{<}(u_2^\varepsilon + u_3^\varepsilon, \tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon)] + u^{\varepsilon, \sharp}$$

with $u^{\varepsilon, \sharp}(t) \in \mathcal{C}^{1+\beta}$. This yields that

$$\|u_3^\varepsilon(t)\|_{1/2+\delta} \lesssim \|u_2^\varepsilon(t) + u_3^\varepsilon(t)\|_\gamma (C_W^\varepsilon + E_W^\varepsilon) + \|u^{\varepsilon, \sharp}(t)\|_{1/2+\delta}, \quad (4.4)$$

and

$$\|u_3^\varepsilon(t)\|_{1-\delta} \lesssim \|u_2^\varepsilon(t) + u_3^\varepsilon(t)\|_\gamma (C_W^\varepsilon + E_W^\varepsilon) + \|u^{\varepsilon, \sharp}(t)\|_{1-\delta}. \quad (4.5)$$

Then u_3^ε solves (4.2) if and only if $u^{\varepsilon, \sharp}$ solves the following equation:

$$u^{\varepsilon, \sharp} = P_t^\varepsilon(\text{Ext} u_0 - u_1^\varepsilon(0)) - \int_0^t P_{t-s}^\varepsilon \left[Q_N[6u_1^\varepsilon \diamond u_2^\varepsilon u_3^\varepsilon + 3u_1^\varepsilon (u_3^\varepsilon)^2 + 3u_1^\varepsilon \diamond (u_2^\varepsilon)^2 + (u_2^\varepsilon + u_3^\varepsilon)^3] \right. \\ \left. + 3P_N[(\pi_{>} + \pi_{0,\diamond})(u_2^\varepsilon + u_3^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})] - 9\varphi^\varepsilon u^\varepsilon \right] ds \\ - 3 \int_0^t P_{t-s}^\varepsilon P_N[\pi_{<}(u_2^\varepsilon + u_3^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})] ds + 3P_N[\pi_{<}(u_2^\varepsilon + u_3^\varepsilon, \tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon)] \\ := P_t^\varepsilon(\text{Ext} \Phi_0^\varepsilon - u_1^\varepsilon(0)) + \int_0^t P_{t-s}^\varepsilon [Q_N \phi_1^{\varepsilon, \sharp} + P_N \phi_2^{\varepsilon, \sharp} + 9\varphi^\varepsilon u^\varepsilon] ds + F^\varepsilon, \quad (4.6)$$

where F^ε represents the last two terms.

First we prove the estimate for $\phi_1^{\varepsilon, \sharp}$.

Proposition 4.2 For $\phi_1^{\varepsilon, \sharp}$ defined above, the following estimate holds:

$$\|Q_N \phi_1^{\varepsilon, \sharp}\|_{-1/2-\delta-2\kappa} \lesssim C(C_W^\varepsilon, E_W^\varepsilon) (1 + \|u^{\varepsilon, \sharp}\|_{1/2+\delta} \|u_3^\varepsilon\|_\gamma + \|u_3^\varepsilon\|_\gamma^2) + \|u_3^\varepsilon\|_\gamma^3.$$

Here the constant we omit is independent of N .

Proof Since

$$\Pi_N[u_3^\varepsilon u_2^\varepsilon u_1^\varepsilon] = P_N[u_3^\varepsilon e_N^{i_1 i_2 i_3} u_2^\varepsilon u_1^\varepsilon],$$

we have for $\delta > 2\kappa$

$$\begin{aligned} & \|\Pi_N[u_3^\varepsilon u_2^\varepsilon u_1^\varepsilon]\|_{-1/2-\delta/2-2\kappa} \lesssim \|u_3^\varepsilon u_2^\varepsilon u_1^\varepsilon e_N^{i_1 i_2 i_3}\|_{-1/2-\delta/2-\kappa} \\ & \lesssim (\|e_N^{i_1 i_2 i_3} u_1^\varepsilon\|_{-1/2-\delta/2-\kappa} \|u_2^\varepsilon\|_{1/2-\delta} + \|\pi_0(u_2^\varepsilon, e_N^{i_1 i_2 i_3} u_1^\varepsilon)\|_{-\delta}) \|u_3^\varepsilon\|_{1/2+\delta} \\ & \lesssim (N^{-\kappa/2} \|u_2^\varepsilon\|_{1/2-\delta} \|u_1^\varepsilon\|_{-1/2-\delta/2} + \|\pi_0(u_2^\varepsilon, e_N^{i_1 i_2 i_3} u_1^\varepsilon)\|_{-\delta}) \|u_3^\varepsilon\|_{1/2+\delta}, \end{aligned}$$

where in the first and last inequalities we used Lemma 3.1.

Using paraproduct one has

$$\begin{aligned} \Pi_N[u_1^\varepsilon (u_3^\varepsilon)^2] &= P_N[u_1^\varepsilon e_N^{i_1 i_2 i_3} (u_3^\varepsilon)^2] \\ &= P_N[\pi_<((u_3^\varepsilon)^2, e_N^{i_1 i_2 i_3} u_1^\varepsilon) + \pi_0((u_3^\varepsilon)^2, e_N^{i_1 i_2 i_3} u_1^\varepsilon) + \pi_>((u_3^\varepsilon)^2, e_N^{i_1 i_2 i_3} u_1^\varepsilon)] \\ &= P_N[\pi_<((u_3^\varepsilon)^2, e_N^{i_1 i_2 i_3} u_1^\varepsilon) + \pi_0(\pi_0(u_3^\varepsilon, u_3^\varepsilon), e_N^{i_1 i_2 i_3} u_1^\varepsilon) \\ &\quad + \pi_>((u_3^\varepsilon)^2, e_N^{i_1 i_2 i_3} u_1^\varepsilon) + 2C(u_3^\varepsilon, u_3^\varepsilon, e_N^{i_1 i_2 i_3} u_1^\varepsilon) + 2u_3^\varepsilon \pi_0(u_3^\varepsilon, e_N^{i_1 i_2 i_3} u_1^\varepsilon)]. \end{aligned}$$

Here $C(u_3^\varepsilon, u_3^\varepsilon, e_N^{i_1 i_2 i_3} u_1^\varepsilon)$ is defined in Lemma 2.3. Then by using Lemmas 2.3 and 3.1 we obtain

$$\|\Pi_N[(u_3^\varepsilon)^2 u_1^\varepsilon]\|_{-1/2-\delta/2-2\kappa} \lesssim N^{-\kappa/2} \|u_3^\varepsilon\|_{1/2+\delta} \|u_3^\varepsilon\|_\gamma \|u_1^\varepsilon\|_{-1/2-\delta/2}. \quad (4.7)$$

Moreover by a similar argument as (4.7) we have

$$\begin{aligned} & \|\Pi_N[(u_2^\varepsilon)^2 \diamond u_1^\varepsilon]\|_{-1/2-\delta/2-2\kappa} \\ & \lesssim N^{-\kappa/2} \|u_2^\varepsilon\|_{1/2-\delta}^2 \|u_1^\varepsilon\|_{-1/2-\delta/2} + \|u_2^\varepsilon\|_{1/2-\delta} \|\pi_{0,\diamond}(u_2^\varepsilon, u_1^\varepsilon e_N^{i_1 i_2 i_3})\|_{-\delta}. \end{aligned}$$

Furthermore Lemma 3.1 implies that

$$\|Q_N[(u_2^\varepsilon + u_3^\varepsilon)^3]\|_{\gamma-\kappa} \lesssim \|u_2^\varepsilon + u_3^\varepsilon\|_\gamma^3.$$

The estimate for the terms containing P_N can be obtained similarly. Hence the result follows from (4.4) and the above estimates. \square

Now we consider $\phi_2^{\varepsilon, \sharp}$. To prove an estimate for $\pi_0(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})$ we have to use paracontrolled ansatz. However, the Fourier cutoff operator P_N does not commute with the paraproduct. Here we follow the technique from [GP15, Lemma 8.16] and prove the following result.

Lemma 4.3 Let $\alpha + \beta + \gamma > 0, \beta + \gamma < 0$, assume that $\alpha \in (0, 1)$, and let $\varphi \in \mathcal{C}^\alpha, \psi \in \mathcal{C}^\beta, \chi \in \mathcal{C}^\gamma$. Define the operator for any $f \in \mathcal{C}^\alpha$

$$A_N^1(\psi, \chi)(f) := \pi_0((I - P_N)\pi_<(f, P_N\psi), \chi),$$

and

$$A_N^2(\psi, \chi)(f) := \pi_0(P_N\pi_<(f, (P_{3N} - P_N)\psi), \chi).$$

Then for all $\eta < 0$

$$\begin{aligned} & \|\pi_0(P_N\pi_<(\varphi, P_{3N}\psi), \chi) - \varphi\pi_0(P_N\psi, \chi)\|_\eta \\ & \lesssim \|\varphi\|_\alpha \|P_N\psi\|_\beta \|\chi\|_\gamma + (\|A_N^1(\psi, \chi)\|_{L(\mathcal{C}^\alpha, \mathcal{C}^\eta)} + \|A_N^2(\psi, \chi)\|_{L(\mathcal{C}^\alpha, \mathcal{C}^\eta)}) \|\varphi\|_\alpha. \end{aligned}$$

Here the constant we omit is independent of N .

Proof We have that

$$\pi_0(P_N \pi_{<}(\varphi, P_{3N} \psi), \chi) = A_N^2(\psi, \chi)(\varphi) + \pi_0(\pi_{<}(\varphi, P_N \psi), \chi) - A_N^1(\psi, \chi)(\varphi).$$

Thus the result follows from Lemma 2.3. \square

Proposition 4.4 For $\phi_2^{\varepsilon, \sharp}$ defined in (4.6), the following estimate holds:

$$\|P_N \phi_2^{\varepsilon, \sharp}\|_{-1/2-2\delta-\kappa} \lesssim C(C_W^\varepsilon, E_W^\varepsilon, A_N, D_N)(1 + \|u_3^\varepsilon\|_\gamma + \|u^{\varepsilon, \sharp}\|_{1+\beta}).$$

with

$$\begin{aligned} A_N := & \|A_N^1(K^\varepsilon + K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})\|_{C_T L(\mathcal{C}^{1-\delta}, \mathcal{C}^{-1/2-2\delta})} \\ & + \|A_N^2(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})\|_{C_T L(\mathcal{C}^{1-\delta}, \mathcal{C}^{-1/2-2\delta})} \end{aligned}$$

and

$$\begin{aligned} D_N := & \sup_{t \in [0, T]} (\|\pi_0((I - P_N) \pi_{<}(u_2^\varepsilon, K^\varepsilon + K_1^\varepsilon), (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})\|_{-\delta} \\ & + \|\pi_0(P_N \pi_{<}(u_2^\varepsilon, (P_{3N} - P_N)(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon)), (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})\|_{-\delta}). \end{aligned}$$

Proof First we consider $\pi_0(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})$. By paracontrolled ansatz we obtain

$$\begin{aligned} & \pi_0(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}) \\ = & -3\pi_0(P_N(\pi_{<}(u_2^\varepsilon + u_3^\varepsilon, P_{3N}(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon)), (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}) + \pi_0(u^{\varepsilon, \sharp}, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}). \end{aligned}$$

Here in the equality we used $P_{3N} \tilde{K}^\varepsilon = \tilde{K}^\varepsilon$. Then by using Lemma 4.3 and $P_N \tilde{K}^\varepsilon = K^\varepsilon$ we obtain that for $\beta > \delta/2$

$$\begin{aligned} & \|\pi_0(u_3^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})\|_{-1/2-2\delta} \\ \lesssim & \|u_2^\varepsilon + u_3^\varepsilon\|_{1/2-\delta} (\|K^\varepsilon + K_1^\varepsilon\|_{1-\delta} \|(u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}\|_{-1-\delta/2} + \|\pi_{0, \diamond}(K^\varepsilon + K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})\|_{-\delta}) \\ & + A_N \|u_3^\varepsilon\|_{1-\delta} + D_N + \|u^{\varepsilon, \sharp}\|_{1+\beta} \|(u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}\|_{-1-\delta/2}. \end{aligned}$$

The estimate for $\pi_{>}(u_2^\varepsilon + u_3^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})$ can be obtained by Lemma 2.2. Thus the result follows by (4.3), (4.4) and (4.5). \square

Remark 4.5 In our case to use the random operator technique, it requires that $u_3^\varepsilon \in \mathcal{C}^{1/2+\beta+\kappa}$. However the best regularity we can obtain for u_2^ε is in $\mathcal{C}^{1/2-\delta}$. Thus for the error terms including u_2^ε we have to calculate it directly which corresponds to D_N .

Estimate for F^ε We now turn to F^ε : Here we divide F^ε into two parts.

$$\begin{aligned} & \|F^\varepsilon\|_{1+\beta} \\ \lesssim & \left\| \int_0^t P_{t-s}^\varepsilon \pi_{<}(u_2^\varepsilon(s) + u_3^\varepsilon(s) - (u_2^\varepsilon(t) + u_3^\varepsilon(t)), (u_1^\varepsilon)^{\diamond, 2}(s) + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}(s)) ds \right\|_{1+\beta} \\ & + \left\| \int_0^t P_{t-s}^\varepsilon \pi_{<}(u_2^\varepsilon(t) + u_3^\varepsilon(t), (u_1^\varepsilon)^{\diamond, 2}(s) + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}(s)) ds - P_N \pi_{<}(u_2^\varepsilon(t) + u_3^\varepsilon(t), \tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon) \right\|_{1+\beta} \\ = & I_1 + I_2. \end{aligned}$$

The estimate for I_2 can be obtained by Lemma 3.3:

$$I_2 \lesssim t^{\frac{\gamma-\beta-\frac{\delta}{2}-\kappa}{2}} \|u_2^\varepsilon(t) + u_3^\varepsilon(t)\|_\gamma (C_W^\varepsilon + E_W^\varepsilon), \quad (4.8)$$

where by the condition on β we have $\frac{\gamma-\beta-\frac{\delta}{2}-\kappa}{2} > 0$.

For I_1 we will use the regularity of $u_2^\varepsilon + u_3^\varepsilon$ with respect to time to control it. Lemmas 2.2 and 3.2 yield that for $5\delta/4 + \beta/2 + \kappa < 1/4$

$$\begin{aligned} I_1 &\lesssim \int_0^t (t-s)^{-1-\frac{\delta/2+\beta+\kappa}{2}} \|(u_1^\varepsilon)^{\diamond,2}(s) + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2}(s)\|_{-1-\delta/2} \|u_2^\varepsilon(t) + u_3^\varepsilon(t) - u_2^\varepsilon(s) - u_3^\varepsilon(s)\|_{\kappa/2} ds \\ &\lesssim (C_W^\varepsilon + E_W^\varepsilon) (C_W^\varepsilon + \int_0^t (t-s)^{-1-\frac{\delta/2+\beta+\kappa}{2}} \|u_3^\varepsilon(t) - u_3^\varepsilon(s)\|_{\kappa/2} ds), \end{aligned}$$

and we note that by Lemmas 3.2 and 3.4 that for $t > s > 0$

$$\begin{aligned} &\|u_3^\varepsilon(t) - u_3^\varepsilon(s)\|_{\kappa/2} \\ &\lesssim \|(P_{\frac{t}{2}}^\varepsilon - P_{\frac{s}{2}}^\varepsilon)(P_{\frac{t}{2}}^\varepsilon + P_{\frac{s}{2}}^\varepsilon)(\text{Ext}\Phi_0^\varepsilon - u_1^\varepsilon(0))\|_{\frac{\kappa}{2}} + \left\| \int_0^s (P_{t-r}^\varepsilon - P_{s-r}^\varepsilon) G^\varepsilon(r) dr \right\|_{\frac{\kappa}{2}} + \left\| \int_s^t P_{t-r}^\varepsilon G^\varepsilon(r) dr \right\|_{\frac{\kappa}{2}} \\ &\lesssim (t-s)^{b_0} s^{-\frac{z+2\kappa+2b_0}{2}} \|\text{Ext}\Phi_0^\varepsilon - u_1^\varepsilon(0)\|_{-z} + (t-s)^b \int_0^s (s-r)^{-\frac{1+\delta+\kappa+2b}{2}} \|G^\varepsilon(r)\|_{-1-\delta} dr \\ &\quad + (t-s)^{b_1} \left(\int_s^t (t-r)^{-\frac{1+\delta+\kappa}{2(1-b_1)}} \|G^\varepsilon(r)\|_{-1-\delta}^{\frac{1}{1-b_1}} dr \right)^{1-b_1}, \end{aligned}$$

where in the last inequality for the third term we used Hölder's inequality. Here $\frac{\delta}{2} + \beta + 2\kappa < 2b_0 < 2 - z - 2\kappa$, $\frac{\delta}{2} + \beta + 2\kappa < 2b < 1 - \kappa - \delta$, $\frac{1}{2}(\frac{\delta}{2} + \beta + 2\kappa) < b_1 < [1 - \frac{3(\gamma+z+\kappa)}{2}] \wedge \frac{1}{2}(1 - \delta - \kappa)$ and

$$G^\varepsilon = Q_N [3u_1^\varepsilon \diamond (u_2^\varepsilon)^2 + 6u_1^\varepsilon \diamond u_2^\varepsilon u_3^\varepsilon + 3u_1^\varepsilon (u_3^\varepsilon)^2 + 3(u_1^\varepsilon)^{\diamond,2} \diamond u_2^\varepsilon + 3(u_1^\varepsilon)^{\diamond,2} \diamond u_3^\varepsilon + (u_2^\varepsilon + u_3^\varepsilon)^3].$$

Moreover, by Propositions 4.2 and 4.3 one has the following estimate

$$\|G^\varepsilon\|_{-1-\delta} \lesssim C(C_W^\varepsilon, E_W^\varepsilon, A_N, D_N) (1 + \|u^{\varepsilon,\sharp}\|_{1/2+\delta} \|u_3^\varepsilon\|_\gamma + \|u_3^\varepsilon\|_\gamma^3 + \|u^{\varepsilon,\sharp}\|_{1+\beta}). \quad (4.9)$$

Thus we obtain that

$$\begin{aligned} I_1 &\lesssim (C_W^\varepsilon + E_W^\varepsilon) \left(C_W^\varepsilon + t^{-\frac{\delta/2+\beta+z}{2}-2\kappa} \|\text{Ext}\Phi_0^\varepsilon - u_1^\varepsilon(0)\|_{-z} \right. \\ &\quad + \int_0^t \int_r^t (t-s)^{-1-\frac{\delta/2+\beta+2\kappa}{2}+b} (s-r)^{-\frac{1+\delta+\kappa+2b}{2}} ds \|G^\varepsilon(r)\|_{-1-\delta} dr \\ &\quad + \left(\int_0^t (t-s)^{-1-\frac{\delta/2+\beta+2\kappa}{2}+b_1} ds \right)^{b_1} \left(\int_0^t \int_0^r (t-s)^{-1-\frac{\delta/2+\beta+2\kappa}{2}+b_1} (t-r)^{-\frac{1+\delta+\kappa}{2(1-b_1)}} \right. \\ &\quad \left. \left. \|G^\varepsilon(r)\|_{-1-\delta}^{\frac{1}{1-b_1}} ds dr \right)^{1-b_1} \right), \end{aligned}$$

where for the last term we used Hölder's inequality. Then by changing variable $s = r + (t - r)\sigma$ for the third term and using (4.9) we have

$$\begin{aligned}
I_1 &\lesssim (C_W^\varepsilon + E_W^\varepsilon) t^{-\frac{\delta/2+\beta+z}{2}-2\kappa} \|\text{Ext}\Phi_0^\varepsilon - u_1^\varepsilon(0)\|_{-z} + C(C_W^\varepsilon, E_W^\varepsilon, A_N, D_N) \\
&\quad + C(C_W^\varepsilon, E_W^\varepsilon, A_N, D_N) \int_0^t (t-r)^{-\frac{1}{2}-\frac{3\delta/2+\beta+3\kappa}{2}} (\|u^{\varepsilon,\sharp}\|_{1/2+\delta} \|u_3^\varepsilon\|_\gamma + \|u^{\varepsilon,\sharp}\|_{1+\beta} + \|u_3^\varepsilon\|_\gamma^3) dr \\
&\quad + C(C_W^\varepsilon, E_W^\varepsilon, A_N, D_N) \left[\int_0^t (t-r)^{-\frac{1+\delta+\kappa}{2(1-b_1)}} (\|u^{\varepsilon,\sharp}\|_{1+\beta} + \|u^{\varepsilon,\sharp}\|_{1/2+\delta} \|u_3^\varepsilon\|_\gamma + \|u_3^\varepsilon\|_\gamma^3)^{\frac{1}{1-b_1}} dr \right]^{1-b_1}.
\end{aligned} \tag{4.10}$$

Combining (4.8) and (4.10) we could control $\|F^\varepsilon\|_{1+\beta}$ by the right hand side of (4.8) and (4.10). Now we also want to estimate $\|F^\varepsilon\|_{1/2+\delta}$ and $\|F^\varepsilon\|_\gamma$. The estimates for these two terms are much easier. We do not need to use Lemma 3.3. We can obtain the following estimates by Lemmas 2.2 and 3.2 directly:

$$\begin{aligned}
&\|F^\varepsilon\|_{1/2+\delta} \\
&\lesssim (C_W^\varepsilon + E_W^\varepsilon) \int_0^t (t-s)^{-\frac{3+3\delta+\kappa}{4}} \|u_3^\varepsilon\|_\gamma ds + C(C_W^\varepsilon, E_W^\varepsilon) + t^{\frac{1-3\delta-\kappa}{4}} \|u_2^\varepsilon(t) + u_3^\varepsilon(t)\|_\gamma (C_W^\varepsilon + E_W^\varepsilon),
\end{aligned} \tag{4.11}$$

and

$$\begin{aligned}
&\|F^\varepsilon\|_\gamma \\
&\lesssim C_W^\varepsilon \int_0^t (t-r)^{-\frac{1+\delta/2+\kappa+\gamma}{2}} \|u_3^\varepsilon\|_\gamma dr + C(C_W^\varepsilon, E_W^\varepsilon) + t^{\frac{2-\delta-2\gamma-\kappa}{4}} \|u_2^\varepsilon(t) + u_3^\varepsilon(t)\|_\gamma (C_W^\varepsilon + E_W^\varepsilon).
\end{aligned} \tag{4.12}$$

Uniform estimate of the solution

Now we introduce the following random time: Define for any $L \geq 1$

$$\tau_L^\varepsilon := \inf\{t \geq 0 : \|u^\varepsilon(t)\|_{-z} \geq L\} \wedge L \quad \rho_L^\varepsilon := \inf\{t \geq 0 : C_W^\varepsilon(t) + E_W^\varepsilon + A_N + D_N \geq L\}.$$

Proposition 4.6 For any $L, L_1 \geq 1$, we have

$$\sup_{t \in [0, \tau_L^\varepsilon \wedge \rho_{L_1}^\varepsilon]} (t^{\frac{3(\gamma+z+\kappa)}{2}} \|u^{\varepsilon,\sharp}\|_{1+\beta} + t^{\frac{1/2+\delta+z+\kappa}{2}} \|u^{\varepsilon,\sharp}\|_{1/2+\delta} + t^{\frac{\gamma+z+\kappa}{2}} \|u^{\varepsilon,\sharp}(t)\|_\gamma) \lesssim C(L, L_1).$$

Moreover before $\tau_L^\varepsilon \wedge \rho_{L_1}^\varepsilon$ one has u_3^ε depends in a Lipschitz continuous way on the data $\text{Ext}\Phi_0^\varepsilon$ and terms in $(C_W^\varepsilon, E_W^\varepsilon, A_N)$. Here we consider u_3^ε with respect to

$$\sup_{t \in [0, \tau_L^\varepsilon \wedge \rho_{L_1}^\varepsilon]} \|u_3^\varepsilon(t)\|_{-z}.$$

Proof By paracontrolled ansatz Lemma 2.2 and (4.3) one then has for $t \in [0, \tau_L^\varepsilon \wedge \rho_{L_1}^\varepsilon]$

$$\|u_3^\varepsilon(t)\|_\gamma \lesssim t^{\delta/4} L_1 \|u_2^\varepsilon(t) + u_3^\varepsilon(t)\|_\gamma + \|u^{\varepsilon,\sharp}(t)\|_\gamma,$$

which shows that for t small enough (depending on L_1)

$$\|u_3^\varepsilon\|_\gamma \lesssim L_1^2 + \|u^{\varepsilon,\sharp}\|_\gamma.$$

Then it follows from Propositions 4.2 4.4 and (4.8) (4.10) that for $\frac{3(\gamma+z+\kappa)}{2} < 1$ and t small enough (depending on L_1)

$$\begin{aligned}
& t^{\frac{3(\gamma+z+\kappa)}{2}} \|u^{\varepsilon, \sharp}(t)\|_{1+\beta} \\
& \lesssim C \|\text{Ext} \Phi_0^\varepsilon - u_1^\varepsilon(0)\|_{-z} + t^{\frac{3(\gamma+z+\kappa)}{2}} C \int_0^t (t-r)^{-\frac{3}{4}-\delta-\frac{\beta}{2}-\kappa} (r^{-\frac{3(\gamma+z+\kappa)}{2}} U^\varepsilon + r^{-\frac{(\gamma+z+\kappa)}{2}-\rho} \|u^{\varepsilon, \sharp}\|_\gamma) dr + C \\
& + C t^{\frac{3(\gamma+z+\kappa)}{2}} \int_0^t (t-r)^{-\frac{1}{2}-\frac{3\delta/2+\beta+3\kappa}{2}} r^{-\frac{3(\gamma+z+\kappa)}{2}} U^\varepsilon(r) dr + t^{\frac{3(\gamma+z+\kappa)}{2(1-b_1)}} \int_0^t (t-r)^{-\frac{1+\delta+\kappa}{2(1-b_1)}} r^{-\frac{3(\gamma+z+\kappa)}{2(1-b_1)}} U^\varepsilon(r)^{\frac{1}{1-b_1}} dr.
\end{aligned} \tag{4.13}$$

Here and in the following $C = C(L_1)$ and

$$U^\varepsilon(r) = r^{3(\gamma+z+\kappa)/2} (\|u^{\varepsilon, \sharp}(r)\|_{1+\beta} + \|u^{\varepsilon, \sharp}(r)\|_{1/2+\delta} \|u^{\varepsilon, \sharp}(r)\|_\gamma + \|u^{\varepsilon, \sharp}(r)\|_\gamma^3).$$

A similar argument as (4.13) and using (4.11) (4.12) one also has that for t small enough (depending on L_1) and $0 < 6\kappa < \frac{3}{2} - 2z - 2\delta - 3\gamma$

$$\begin{aligned}
& t^{\frac{1/2+\delta+z+\kappa}{2}} \|u^{\varepsilon, \sharp}(t)\|_{1/2+\delta} \\
& \lesssim \|\text{Ext} \Phi_0^\varepsilon - u_1^\varepsilon(0)\|_{-z} + t^{\frac{1/2+\delta+z+\kappa}{2}} C \int_0^t (t-s)^{-\frac{1+3\delta+2\kappa}{2}} (s^{-\frac{3(\gamma+z+\kappa)}{2}} U^\varepsilon + s^{-\frac{(\gamma+z+\kappa)}{2}-\rho} \|u^{\varepsilon, \sharp}\|_\gamma) ds \\
& + C + C t^{\frac{1/2+\delta+z+\kappa}{2}} \int_0^t (t-r)^{-\frac{3+3\delta+\kappa}{4}} r^{-\frac{\gamma+z+\kappa}{2}} r^{\frac{\gamma+z+\kappa}{2}} \|u^{\varepsilon, \sharp}\|_\gamma dr,
\end{aligned} \tag{4.14}$$

and

$$\begin{aligned}
& t^{\frac{\gamma+z+\kappa}{2}} \|u^{\varepsilon, \sharp}(t)\|_\gamma \\
& \lesssim \|\text{Ext} \Phi_0^\varepsilon - u_1^\varepsilon(0)\|_{-z} + t^{\frac{\gamma+z+\kappa}{2}} C \int_0^t (t-s)^{-\frac{1}{4}-\delta-\frac{\gamma}{2}-\kappa} (s^{-\frac{3(\gamma+z+\kappa)}{2}} U^\varepsilon + s^{-\frac{(\gamma+z+\kappa)}{2}-\rho} \|u^{\varepsilon, \sharp}\|_\gamma) ds \\
& + C + C t^{\frac{\gamma+z+\kappa}{2}} \int_0^t (t-r)^{-\frac{1+\delta/2+\gamma+\kappa}{2}} r^{-\frac{(\gamma+z+\kappa)}{2}} r^{\frac{\gamma+z+\kappa}{2}} \|u^{\varepsilon, \sharp}\|_\gamma dr.
\end{aligned} \tag{4.15}$$

Since $\frac{1/2+\delta+\kappa+z}{2} \leq \gamma + z + \kappa$, combining (4.13-4.15) we get that by Bihari's inequality there exists some T_0 (depending on L_1) such that

$$\sup_{t \in [0, T_0]} (t^{\frac{3(\gamma+z+\kappa)}{2}} \|u^{\varepsilon, \sharp}\|_{1+\beta} + t^{\frac{1/2+\delta+z+\kappa}{2}} \|u^{\varepsilon, \sharp}\|_{1/2+\delta} + t^{\frac{\gamma+z+\kappa}{2}} \|u^{\varepsilon, \sharp}(t)\|_\gamma) \lesssim C(L, L_1),$$

which combining with Propositions 4.2 and 4.4 implies that

$$\sup_{t \in [0, T_0]} t^{3(\gamma+z+\kappa)/2} \|Q_N \phi_1^{\varepsilon, \sharp} + P_N \phi_2^{\varepsilon, \sharp}\|_{-1/2-2\delta-\kappa} \lesssim C(L, L_1). \tag{4.16}$$

Moreover by paracontrolled ansatz we also obtain

$$\|u_3^\varepsilon\|_{-z} \lesssim t^{\delta/4} \|u_2^\varepsilon + u_3^\varepsilon\|_{-z} L_1 + \|u^{\varepsilon, \sharp}\|_{-z},$$

which combining with (4.16) implies that for t small enough and $t \in [0, T_0]$

$$\begin{aligned}
\|u_3^\varepsilon(t)\|_{-z} &\lesssim C + \|u^{\varepsilon, \#}(t)\|_{-z} \\
&\lesssim C + \|\text{Ext}\Phi_0^\varepsilon - u_1^\varepsilon(0)\|_{-z} + \|F^\varepsilon(t)\|_{-z} + \int_0^t (t-s)^{-\frac{1/2+2\delta+3\kappa-z}{2}} s^{-\rho} \|u^\varepsilon\|_{-1/2-\delta} ds \\
&\quad + \int_0^t (t-s)^{-\frac{1/2+2\delta+3\kappa-z}{2}} s^{-\frac{3(\gamma+z+\kappa)}{2}} s^{\frac{3(\gamma+z+\kappa)}{2}} \|Q_N \phi_1^{\varepsilon, \#} + P_N \phi_2^{\varepsilon, \#}\|_{-1/2-2\delta-\kappa} ds \\
&\lesssim C(L, L_1) + t^{\frac{1}{2}-\frac{\delta}{4}-\kappa} \|u_3^\varepsilon(t)\|_{-z} C.
\end{aligned}$$

Here in the last inequality we used

$$\begin{aligned}
&\|F^\varepsilon(t)\|_{-z} \\
&\lesssim C \int_0^t (t-s)^{-\frac{1+\delta/2+\kappa-z}{2}} s^{-\frac{\gamma+\kappa+z}{2}} ds \sup_{s \in [0, t]} s^{\frac{\gamma+\kappa+z}{2}} \|u_2^\varepsilon + u_3^\varepsilon\|_\gamma + t^{\frac{1}{2}-\frac{\delta}{4}-\kappa} \|u_2^\varepsilon(t) + u_3^\varepsilon(t)\|_{-z} C.
\end{aligned}$$

Hence before T_0 one has u_3^ε depends in a Lipschitz continuous way on the data $\text{Ext}\Phi_0^\varepsilon$ and terms in $(C_W^\varepsilon, E_W^\varepsilon, A_N)$. Furthermore we can extend the time to $\tau_L^\varepsilon \wedge \rho_{L_1}^\varepsilon$ as we did in [ZZ14]. \square

5 Proof of main result

In [CC13] it is obtained that the solution to (1.1) can be obtained as limit of solutions $\bar{\Phi}^\varepsilon$ to the following equation:

$$d\bar{\Phi}^\varepsilon = \Delta \bar{\Phi}^\varepsilon dt + P_N dW - (\bar{\Phi}^\varepsilon)^3 dt + (3\bar{C}_0^\varepsilon - 9\bar{C}_1^\varepsilon) \bar{\Phi}^\varepsilon dt,$$

$$\bar{\Phi}^\varepsilon(0) = \Phi_0.$$

Here \bar{C}_0^ε and \bar{C}_1^ε are defined in Section 6.1. For $L \geq 0$ define $\tau_L := \inf\{t \geq 0 : \|\Phi(t)\|_{-z} \geq L\} \wedge L$ and then τ_L increases to the explosion time τ . Moreover define $\bar{\tau}_L^\varepsilon := \inf\{t \geq 0 : \|\bar{\Phi}^\varepsilon(t)\|_{-z} \geq L\} \wedge L$ and $\bar{\rho}_L^\varepsilon := \inf\{t \geq 0 : \bar{C}_W^\varepsilon(t) \geq L\}$ with \bar{C}_W^ε defined similarly as C_W^ε . A similar argument as above implies that

$$\sup_{t \in [0, \tau_L \wedge \bar{\rho}_{L_3}^\varepsilon \wedge \bar{\tau}_{L_4}^\varepsilon]} \|\bar{\Phi}^\varepsilon(t) - \Phi(t)\|_{-z} \xrightarrow{P} 0. \quad (5.1)$$

Here Φ is the solution to (1.2). Define

$$\begin{aligned}
\delta C_W^\varepsilon &:= \sup_{t \in [0, T]} (\|u_1^\varepsilon - \bar{u}_1^\varepsilon\|_{-1/2-\delta/2} + \|(u_1^\varepsilon)^{\diamond, 2} - (\bar{u}_1^\varepsilon)^{\diamond, 2}\|_{-1-\delta/2} + \|u_2^\varepsilon - \bar{u}_2^\varepsilon\|_{1/2-\delta} \\
&\quad + \|\pi_0(u_2^\varepsilon, u_1^\varepsilon) - \pi_0(\bar{u}_2^\varepsilon, \bar{u}_1^\varepsilon)\|_{-\delta} + \|\pi_{0, \diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond, 2}) - \pi_{0, \diamond}(\bar{u}_2^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond, 2})\|_{-1/2-\delta} \\
&\quad + \|\pi_{0, \diamond}(K^\varepsilon, (u_1^\varepsilon)^{\diamond, 2}) - \pi_{0, \diamond}(\bar{K}^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond, 2})\|_{-\delta} + \|u_2^\varepsilon - \bar{u}_2^\varepsilon\|_{C_T^{1/4-\delta-\kappa/2} C_{\kappa/2}}).
\end{aligned}$$

Here $\bar{u}_1^\varepsilon, \bar{u}_2^\varepsilon, \bar{u}_3^\varepsilon$ and associated terms are defined similarly as $u_1^\varepsilon, u_2^\varepsilon, u_3^\varepsilon$ and associated terms respectively. In Section 6 we will prove that $\delta C_W^\varepsilon \xrightarrow{P} 0, E_W^\varepsilon \xrightarrow{P} 0, A_N \xrightarrow{P} 0$ and $D_N \xrightarrow{P} 0$ as $\varepsilon \rightarrow 0$. Then by a similar argument as Section 4 we have

$$\sup_{t \in [0, \tau_L \wedge \tau_{L_1}^\varepsilon \wedge \rho_{L_2}^\varepsilon \wedge \bar{\rho}_{L_3}^\varepsilon \wedge \bar{\tau}_{L_4}^\varepsilon]} \|u^\varepsilon(t) - \bar{\Phi}^\varepsilon(t)\|_{-z} \xrightarrow{P} 0, \quad \varepsilon \rightarrow 0. \quad (5.2)$$

Here $E_W^\varepsilon, A_N, D_N$ appear as error terms for lattice approximations. Then (5.1) and (5.2) implies that

$$\sup_{t \in [0, \tau_L \wedge \tau_{L_1}^\varepsilon \wedge \rho_{L_2}^\varepsilon \wedge \bar{\rho}_{L_3}^\varepsilon \wedge \bar{\tau}_{L_4}^\varepsilon]} \|u^\varepsilon(t) - \Phi(t)\|_{-z} \xrightarrow{P} 0, \quad \varepsilon \rightarrow 0. \quad (5.3)$$

Moreover we have the following estimates:

$$\begin{aligned} & P\left(\sup_{t \in [0, \tau_L]} \|u^\varepsilon - \Phi\|_{-z} > \epsilon\right) \\ & \leq P\left(\sup_{t \in [0, \tau_L \wedge \tau_{L_1}^\varepsilon \wedge \rho_{L_2}^\varepsilon \wedge \bar{\rho}_{L_3}^\varepsilon \wedge \bar{\tau}_{L_4}^\varepsilon]} \|u^\varepsilon - \Phi\|_{-z} > \epsilon\right) + P(\tau_L \wedge \rho_{L_2}^\varepsilon \wedge \bar{\rho}_{L_3}^\varepsilon \wedge \bar{\tau}_{L_4}^\varepsilon > \tau_{L_1}^\varepsilon) \\ & \quad + P(\tau_L \wedge \bar{\rho}_{L_3}^\varepsilon > \bar{\tau}_{L_4}^\varepsilon) + P(\tau_L > \rho_{L_2}^\varepsilon) + P(\tau_L > \bar{\rho}_{L_3}^\varepsilon). \end{aligned}$$

The first term goes to zero as $\varepsilon \rightarrow 0$ by (5.3). Also for $L_1 > L + \epsilon$

$$P(\tau_L \wedge \rho_{L_2}^\varepsilon \wedge \bar{\rho}_{L_3}^\varepsilon \wedge \bar{\tau}_{L_4}^\varepsilon > \tau_{L_1}^\varepsilon) \leq P\left(\sup_{t \in [0, \tau_L \wedge \tau_{L_1}^\varepsilon \wedge \rho_{L_2}^\varepsilon \wedge \bar{\rho}_{L_3}^\varepsilon \wedge \bar{\tau}_{L_4}^\varepsilon]} \|u^\varepsilon - \Phi\|_{-z} > \epsilon\right),$$

which goes to zero as $\varepsilon \rightarrow 0$ by (5.3). Moreover for $L_4 > L + \epsilon$ we have

$$P(\tau_L \wedge \bar{\rho}_{L_3}^\varepsilon > \bar{\tau}_{L_4}^\varepsilon) \leq P\left(\sup_{t \in [0, \tau_L \wedge \bar{\rho}_{L_3}^\varepsilon \wedge \bar{\tau}_{L_4}^\varepsilon]} \|\bar{\Phi}^\varepsilon - \Phi\|_{-z} > \epsilon\right)$$

which goes to zero by (5.1). The last two terms go to zero uniformly over $\varepsilon \in (0, 1)$ as L_2, L_3 go to ∞ . Thus the result follows. \square

6 Stochastic convergence

To simplify the arguments below, we assume that $\mathcal{F}W(0) = 0$ and restrict ourselves to the flow of $\int_{\mathbb{T}^3} u(x) dx = 0$. We follow the notations from [GP15, Section 9]. We represent the white noise in terms of its spatial Fourier transform. More precisely, let $E = \mathbb{Z}^3 \setminus \{0\}$ and let $W(s, k) = \langle W(s), e_k \rangle$ for $e_k(x) = 2^{-3/2} e^{i\pi x \cdot k}$, $x \in \mathbb{T}^3$. Then

$$u_1^\varepsilon(t, x) = \int_{\mathbb{R} \times E} e_k(x) P_{t-s}^\varepsilon(k) W(d\eta), \quad \bar{u}_1^\varepsilon(t, x) = \int_{\mathbb{R} \times E} e_k(x) \bar{P}_{t-s}^\varepsilon(k) W(d\eta),$$

where $\eta_a = (s_a, k_a)$, $s_{-a} = s_a$, $k_{-a} = -k_a$ and the measure $d\eta_a = ds_a dk_a$ is the product of the Lebesgue measure ds_a on \mathbb{R} and of the counting measure dk_a on E and $p_t^\varepsilon(k) = e^{-|k|^2 f(\varepsilon k) t} 1_{\{t \geq 0\}}$, $\bar{P}_t^\varepsilon(k) = p_t^\varepsilon(k) 1_{\{|k|_\infty \leq N\}}$, $P_t^\varepsilon(k) = p_t^\varepsilon(k) 1_{\{|k|_\infty \leq N\}}$. Moreover,

$$\int P_{t-s}^\varepsilon(k) P_{\sigma-s}^\varepsilon(k) ds = \frac{e^{-|k|^2 f(\varepsilon k) |t-\sigma|} 1_{\{|k|_\infty \leq N\}}}{2|k|^2 f(\varepsilon k)} := V_{t-\sigma}^\varepsilon(k), \quad (6.1)$$

and

$$\int \bar{P}_{t-s}^\varepsilon(k) \bar{P}_{\sigma-s}^\varepsilon(k) ds = \frac{e^{-|k|^2 |t-\sigma|} 1_{\{|k|_\infty \leq N\}}}{2|k|^2} := \bar{V}_{t-\sigma}^\varepsilon(k). \quad (6.2)$$

In this section we will prove that $\delta C_W^\varepsilon \rightarrow 0$, $E_W^\varepsilon \rightarrow 0$, $A_N \rightarrow 0$, $D_N \rightarrow 0$ in probability as $\varepsilon \rightarrow 0$.

Now we introduce the following notations: $k_{[1\dots n]} = \sum_{i=1}^n k_i$, $\tilde{k}^{i_1 i_2 i_3} = (k^j - i_j(2N+1))_{j=1,2,3}$ for $i_j = 1, 0, -1$ and $\sum_{j=1}^3 i_j^2 \neq 0$. In the following we always omit the superscript of \tilde{k} if there's no confusion. Denote by

$$\int_{(\mathbb{R} \times E)^n} f(\eta_{1\dots n}) W(d\eta_{1\dots n})$$

a generic element of the n -th chaos of W . By [GP15, Section 9.2] We know that

$$E[|\int_{(\mathbb{R} \times E)^n} f(\eta_{1\dots n}) W(d\eta_{1\dots n})|^2] \leq (n!) \int_{(\mathbb{R} \times E)^n} |f(\eta_{1\dots n})|^2 d\eta_{1\dots n},$$

such that for bounding the variance of the chaos it is enough to bound the L^2 norm of the unsymmetrized kernels. To obtain the results we first recall the following lemma from [ZZ14] for our later use:

Lemma 6.1 ([ZZ14, Lemma 3.10]) Let $0 < l, m < d, l + m - d > 0$. Then we have

$$\sum_{k_1, k_2 \in \mathbb{Z}^d \setminus \{0\}, k_1 + k_2 = k} \frac{1}{|k_1|^l |k_2|^m} \lesssim \frac{1}{|k|^{l+m-d}}.$$

By a similar argument as the proof of [ZZ14, Lemma 3.11] we have the following results.

Lemma 6.2 For any $0 < \kappa < 1, i \geq 0, t \geq 0, k_1, k_2 \in E$ we have

$$|e^{-|k_{[12]}|^2 t} \theta(2^{-i} k_{[12]}) - e^{-|k_2|^2 t} \theta(2^{-i} k_2)| \lesssim |k_1|^\kappa 2^{-i\kappa}.$$

Lemma 6.3 For any $0 < \kappa < 1, i \geq 0, t \geq 0$ we have for $k_1, k_2 \in E$ with $|k_{[12]}|_\infty \leq N, |k_2|_\infty \leq N$

$$|e^{-|k_{[12]}|^2 t f(\varepsilon k_{[12]})} \theta(2^{-i} k_{[12]}) - e^{-|k_2|^2 t f(\varepsilon k_2)} \theta(2^{-i} k_2)| \lesssim |k_1|^\kappa 2^{-i\kappa}.$$

Now we prove the following estimate for the approximating operators:

Lemma 6.4 For any $0 < \kappa < 1$ and $t > 0, k \in E, \varepsilon > 0$

$$(1) |p_t^\varepsilon(k) - p_t(k)| \lesssim e^{-|k|^2 \bar{c}_f t} |\varepsilon k|^\kappa, \quad |P_t^\varepsilon(k) - p_t(k)| \lesssim e^{-|k|^2 \bar{c}_f t} |\varepsilon k|^\kappa;$$

$$(2) |P_t^\varepsilon(k) - \bar{P}_t^\varepsilon(k)| \lesssim e^{-|k|^2 \bar{c}_f t} |\varepsilon k|^\kappa, \quad |V_t^\varepsilon(k) - \bar{V}_t^\varepsilon(k)| \lesssim \frac{e^{-|k|^2 \bar{c}_f t} |\varepsilon k|^\kappa}{|k|^2}.$$

Here $\bar{c}_f = c_f \wedge 1$, $c_f = \min\{f(x) : |x| \leq 1.8\}$.

Proof The results follow from $|f(\varepsilon k) - 1| \lesssim |\varepsilon k|^\kappa$ and

$$|e^{-|k|^2 t f(\varepsilon k)} - e^{-|k|^2 t}| \lesssim e^{-|k|^2 \bar{c}_f t} (1 \wedge t^\kappa |f(\varepsilon k) - 1| |k|^{2\kappa}) \lesssim e^{-|k|^2 \bar{c}_f t} |\varepsilon k|^\kappa.$$

□

We prove the following two lemmas for dealing with the error terms.

Lemma 6.5 For every $q \geq 0, 0 < r < 3$,

$$\int_E \theta(2^{-q}\tilde{k})^2 \frac{1}{|k|^r} dk \lesssim 2^{(3-r)q}, \quad \int_E \theta(2^{-q}\tilde{k})^2 \frac{1}{|k|^r} dk \lesssim 2^{(3-r)q}.$$

Proof We consider the first one, the second can be obtained by a similar argument. We have

$$\int \theta(2^{-q}\tilde{k})^2 \frac{1}{|k|^r} dk \lesssim \int 1_{|k| \leq 2^q} \theta(2^{-q}\tilde{k})^2 \frac{1}{|k|^r} dk + \int 1_{|k| > 2^q} \theta(2^{-q}\tilde{k})^2 \frac{1}{|k|^r} dk \lesssim 2^{(3-r)q}$$

Here in the last inequality we used that the cardinality of the k with $\theta(2^{-q}\tilde{k}) \neq 0$ is of order 2^{3q} . \square

Lemma 6.6 For every $q \geq 0, 0 < r < 3$,

$$\int \theta(2^{-q}\tilde{k})^2 \frac{1}{|k|^r} dk \lesssim \varepsilon^\kappa 2^{(3-r+\kappa)q}.$$

Here $\kappa > 0$ is small enough.

Proof We have

$$\int \theta(2^{-q}\tilde{k})^2 \frac{1}{|k|^r} dk \lesssim \int 1_{|k| \leq N} \theta(2^{-q}\tilde{k})^2 \frac{1}{|k|^r} dk + \varepsilon^\kappa \int 1_{|k| \geq N} \theta(2^{-q}\tilde{k})^2 \frac{1}{|k|^{r-\kappa}} dk \lesssim \varepsilon^\kappa 2^{(3-r+\kappa)q},$$

where in the last inequality we used that $|k| \leq N \simeq |\tilde{k}| \simeq 2^q$ and Lemma 6.5. \square

6.1 Convergence for renormalisation terms

In this subsection we prove $\delta C_W^\varepsilon \rightarrow 0$ in probability as $\varepsilon \rightarrow 0$.

Convergence for $u_1^\varepsilon - \bar{u}_1^\varepsilon$

In this part we consider the convergence of $u_1^\varepsilon - \bar{u}_1^\varepsilon$.

$$\begin{aligned} & E|\Delta_q[u_1^\varepsilon(t) - \bar{u}_1^\varepsilon(t)]|^2 \\ & \lesssim \int_{\mathbb{R} \times E} \theta(2^{-q}k)^2 |e_k(P_{t-s}^\varepsilon(k) - \bar{P}_{t-s}^\varepsilon(k))|^2 d\eta \lesssim \int \theta(2^{-q}k)^2 (\varepsilon|k|)^\kappa |k|^{-2} dk \lesssim \varepsilon^\kappa 2^{q(\kappa+1)}. \end{aligned}$$

Here $\kappa > 0$ is small enough and in the second inequality we used Lemma 6.4. Similarly by using

$$|1 - e^{-|t_2-t_1|f(\varepsilon k)|k|^2}| \lesssim |t_1 - t_2|^\kappa |k|^{2\kappa},$$

we get desired estimates for $E|\Delta_q[(u_1^\varepsilon(t_2) - \bar{u}_1^\varepsilon(t_2)) - (u_1^\varepsilon(t_1) - \bar{u}_1^\varepsilon(t_1))]|$, which combining with Gaussian hypercontractivity implies that for $p > 1, \epsilon > 0$ small enough

$$\begin{aligned} & E[\|(u_1^\varepsilon(t_2) - \bar{u}_1^\varepsilon(t_2)) - (u_1^\varepsilon(t_1) - \bar{u}_1^\varepsilon(t_1))\|_{B_{p,p}^{-1/2-\kappa-\epsilon}}^p] \\ & \lesssim \varepsilon^{p\kappa/2} |t_2 - t_1|^{\kappa p/4}, \end{aligned}$$

Then by Lemma 2.1 we obtain that for $\delta > 0, p > 1$, $u_1^\varepsilon - \bar{u}_1^\varepsilon \rightarrow 0$ in $L^p(\Omega; C_T \mathcal{C}^{-1/2-\delta/2})$ as $\varepsilon \rightarrow 0$.

Convergence for $u_1^\varepsilon \diamond u_1^\varepsilon - \bar{u}_1^\varepsilon \diamond \bar{u}_1^\varepsilon$

In this part we consider the convergence of $u_1^\varepsilon \diamond u_1^\varepsilon$. Recall that $u_1^\varepsilon \diamond u_1^\varepsilon = u_1^\varepsilon u_1^\varepsilon - C_0^\varepsilon$ and $\bar{u}_1^\varepsilon \diamond \bar{u}_1^\varepsilon = \bar{u}_1^\varepsilon \bar{u}_1^\varepsilon - \bar{C}_0^\varepsilon$.

Take

$$C_0^\varepsilon = 2^{-3} \int_E \frac{1_{\{|k|_\infty \leq N\}}}{2|k|^2 f(\varepsilon k)} dk, \quad \bar{C}_0^\varepsilon = 2^{-3} \int \frac{1_{\{|k|_\infty \leq N\}}}{2|k|^2} dk. \quad (6.3)$$

Then we have

$$\begin{aligned} & E|\Delta_q[u_1^\varepsilon \diamond u_1^\varepsilon(t) - \bar{u}_1^\varepsilon \diamond \bar{u}_1^\varepsilon(t)]|^2 \\ & \lesssim \int_{(\mathbb{R} \times E)^2} \theta(2^{-q}k_{[12]})^2 |(P_{t-s_1}^\varepsilon(k_1)P_{t-s_2}^\varepsilon(k_2) - \bar{P}_{t-s_1}^\varepsilon(k_1)\bar{P}_{t-s_2}^\varepsilon(k_2))|^2 d\eta_{12} \\ & \lesssim \varepsilon^\kappa \int \theta(2^{-q}k_{[12]})^2 \frac{|k_1|^\kappa + |k_2|^\kappa}{|k_1|^2 |k_2|^2} dk_{12} \lesssim \varepsilon^\kappa 2^{(\kappa+2)q}. \end{aligned}$$

Here $\kappa > 0$ is small enough and in the second inequality we used Lemma 6.4 and in the last inequality we used Lemma 6.1. Then by Gaussian hypercontractivity and Lemma 2.1 we obtain that for $\delta > 0, p > 1$, $u_1^\varepsilon \diamond u_1^\varepsilon - \bar{u}_1^\varepsilon \diamond \bar{u}_1^\varepsilon \rightarrow 0$ in $L^p(\Omega; C_T \mathcal{C}^{-1-\delta})$ as $\varepsilon \rightarrow 0$.

Convergence for $u_2^\varepsilon - \bar{u}_2^\varepsilon$

In this part we consider the convergence of u_2^ε . Recall that

$$\bar{u}_2^\varepsilon(t) - u_2^\varepsilon(t) = I_t^3 - \bar{I}_t^3 + J_t^3.$$

Here

$$I_t^3 = 2^{-3} \int_{(\mathbb{R} \times E)^3} e_{k_{[123]}} \int_0^t P_{t-\sigma}^\varepsilon(k_{[123]}) P_{\sigma-s_1}^\varepsilon(k_1) P_{\sigma-s_2}^\varepsilon(k_2) P_{\sigma-s_3}^\varepsilon(k_3) d\sigma W(d\eta_{123}),$$

and \bar{I}_t^3 is defined similarly as I_t^3 with $P_{t-\sigma}^\varepsilon(k_{[123]})$ replaced by $p_{t-\sigma}(k_{[123]})$ and with other P^ε replaced by \bar{P}^ε and J_t^3 is defined similarly as I_t^3 with $e_{k_{[123]}}$, $k_{[123]}$ replaced by $e_{\tilde{k}_{[123]}}$, $\tilde{k}_{[123]}$. By Lemma 6.4 and a straightforward calculation we obtain that

$$\begin{aligned} & E|\Delta_q(I_t^3 - \bar{I}_t^3)|^2 \\ & \lesssim \int_{(\mathbb{R} \times E)^2} \theta(2^{-q}k_{[123]})^2 \left| \int_0^t \left(P_{t-\sigma}^\varepsilon(k_{[123]}) \Pi_{i=1}^3 P_{\sigma-s_i}^\varepsilon(k_i) - p_{t-\sigma}(k_{[123]}) \Pi_{i=1}^3 \bar{P}_{\sigma-s_i}^\varepsilon(k_i) \right) d\sigma \right|^2 d\eta_{123} \\ & \lesssim \int \theta(2^{-q}k_{[123]}) \frac{\varepsilon^\kappa \sum_{i=1}^3 |k_i|^\kappa + |k_{[123]}|^\kappa \varepsilon^\kappa}{|k_1|^2 |k_2|^2 |k_3|^2 [|k_1|^2 + |k_2|^2 + |k_3|^2] |k_{[123]}|^2} dk_{123} \\ & \lesssim \int_E \theta(2^{-q}k) \frac{\varepsilon^\kappa}{|k|^{4-\kappa}} dk \lesssim \varepsilon^\kappa 2^{q(-1+\kappa)}, \end{aligned}$$

where we used Young's inequality and Lemma 6.1 in the second inequality. Similar calculations also imply that

$$\begin{aligned} E|\Delta_q J_t^3|^2 & \lesssim \int \theta(2^{-q}\tilde{k}_{[123]}) \frac{1_{\{|k_{123}| > N, 2^q \lesssim N\}}}{|k_1|^2 |k_2|^2 |k_3|^2 [|k_1|^2 + |k_2|^2 + |k_3|^2] |\tilde{k}_{[123]}|^2} dk_{123} \\ & \lesssim \int_E \theta(2^{-q}\tilde{k}) \frac{\varepsilon^\kappa 1_{\{|k| > N, 2^q \lesssim N\}}}{|k|^{2-2\kappa} |\tilde{k}|^2} dk \lesssim \varepsilon^\kappa 2^{q(-1+\kappa)}, \end{aligned}$$

where we used Lemma 6.5 in the last inequality. By a similar argument as above we also obtain that for $\delta > 0, p > 1$, $u_2^\varepsilon - \bar{u}_2^\varepsilon \rightarrow 0$ in $L^p(\Omega; C_T \mathcal{C}^{1/2-\delta/2})$. Similarly we obtain $u_2^\varepsilon - \bar{u}_2^\varepsilon \rightarrow 0$ in $L^p(\Omega; C^{1/4-\delta-\kappa/2}([0, T], \mathcal{C}^{\kappa/2}))$.

Convergence for $\pi_{0,\diamond}(K^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) - \pi_{0,\diamond}(\bar{K}^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond,2})$

In this part we focus on $\pi_0(K^\varepsilon, (u_1^\varepsilon)^{\diamond,2})$ and prove that $\pi_{0,\diamond}(K^\varepsilon, (u_1^\varepsilon)^{\diamond,2}) - \pi_{0,\diamond}(\bar{K}^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond,2})$ in $C_T \mathcal{C}^{-\delta}$ for every $\delta > 0$. Now we have the following identity: for $t \in [0, T]$,

$$\pi_0(K^\varepsilon, (u_1^\varepsilon)^{\diamond,2})(t) - \pi_0(\bar{K}^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond,2})(t) = I_t^1 + 4I_t^2 + 2I_t^3 - [\bar{I}_t^1 + 4\bar{I}_t^2 + 2\bar{I}_t^3].$$

Here

$$\begin{aligned} I_t^1 &= 2^{-\frac{9}{2}} \int e_{k_{[1234]}} \psi_0(k_{[12]}, k_{[34]}) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[12]}) P_{\sigma-s_1}^\varepsilon(k_1) P_{\sigma-s_2}^\varepsilon(k_2) P_{t-s_3}^\varepsilon(k_3) P_{t-s_4}^\varepsilon(k_4) W(d\eta_{1234}), \\ I_t^2 &= 2^{-\frac{9}{2}} \int \int e_{k_{[23]}} \psi_0(k_{[12]}, k_3 - k_1) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[12]}) P_{\sigma-s_2}^\varepsilon(k_2) P_{t-s_3}^\varepsilon(k_3) V_{t-\sigma}^\varepsilon(k_1) dk_1 W(d\eta_{23}), \\ I_t^3 &= 2^{-6} \int_{E^2} \int_0^t d\sigma V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) P_{t-\sigma}^\varepsilon(k_{[12]}) dk_{12}, \end{aligned}$$

and for $i = 1, 2, 3$, \bar{I}_t^i is defined similarly with $P_{t-\sigma}^\varepsilon(k_{[12]})$ replaced by $p_{t-\sigma}(k_{[12]})$ and other $P^\varepsilon, V^\varepsilon$ replaced by $\bar{P}^\varepsilon, \bar{V}^\varepsilon$ respectively. In fact, choose

$$C_{11}^\varepsilon = 2^{-5} \int \int_{-\infty}^t d\sigma V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) P_{t-\sigma}^\varepsilon(k_{[12]}) dk_{12} \quad (6.4)$$

and \bar{C}_{11}^ε is defined with each $P^\varepsilon, V^\varepsilon$ replaced by p, \bar{V}^ε respectively. Choose $\varphi_1^\varepsilon(t) = 2I_t^3 - C_{11}^\varepsilon$ and $\bar{\varphi}_1^\varepsilon(t) = 2\bar{I}_t^3 - \bar{C}_{11}^\varepsilon$ and $\varphi_1(t) = -2^{-7} \int \frac{e^{-t(|k_1|^2 + |k_2|^2 + |k_{[12]}|^2)}}{|k_1|^2 |k_2|^2 (|k_1|^2 + |k_2|^2 + |k_{[12]}|^2)} dk_{12}$. Then we can easily obtain that

$$\sup_{t \in [0, T]} t^\rho |\varphi_1^\varepsilon - \varphi_1| \lesssim \varepsilon^\kappa, \quad \sup_{t \in [0, T]} t^\rho |\bar{\varphi}_1^\varepsilon - \varphi_1| \lesssim \varepsilon^\kappa,$$

for every $\rho > 0, 0 < \kappa < 2\rho$.

Term in the second chaos: Now we consider I_t^2 and by Lemma 6.4 and (6.1), (6.2) we have the following calculations:

$$\begin{aligned} & E|\Delta_q(I_t^2 - \bar{I}_t^2)|^2 \\ & \lesssim \int \psi_0(k_{[12]}, k_3 - k_1) \psi_0(k_{[24]}, k_3 - k_4) \theta(2^{-q} k_{[23]})^2 \\ & \quad \frac{|\varepsilon k_{[12]}|^{\kappa/2} |\varepsilon k_{[24]}|^{\kappa/2} + |\varepsilon k_1|^{\kappa/2} |\varepsilon k_4|^{\kappa/2} + |\varepsilon k_2|^\kappa + |\varepsilon k_3|^\kappa}{|k_1|^2 |k_2|^2 |k_3|^2 |k_4|^2 (|k_1|^2 + |k_{[12]}|^2) (|k_4|^2 + |k_{[24]}|^2)} dk_{1234} \\ & \lesssim \varepsilon^\kappa \int \theta(2^{-q} k_{[23]})^2 \frac{2^{-2q+2\kappa}}{|k_2|^{2-\kappa} |k_3|^2} dk_{23} \\ & \lesssim \varepsilon^\kappa 2^{q^3 \kappa}, \end{aligned}$$

with $\kappa > 0$ small enough. Here we used $|k_{[i2]}| \gtrsim 2^q$ on the support of $\psi_0(k_{[i2]}, k_3 - k_i) \theta(2^{-q} k_{[23]})$ for $i = 1, 4$ in the second inequality and Lemma 6.1 in the last inequality.

Terms in the fourth chaos: Now for I_t^1 by (6.1), (6.2) and Lemma 6.4 we have the following calculations:

$$\begin{aligned}
& E[|\Delta_q(I_t^1 - \bar{I}_t^1)|^2] \\
& \lesssim \varepsilon^\kappa \int \theta(2^{-q}k_{[1234]})^2 \frac{\theta(2^{-q}k)^2 \psi_0(k_{[12]}, k_{[34]})}{|k_1|^2 |k_2|^2 |k_3|^2 |k_4|^2 |k_{[12]}|^4} (|k_{[12]}|^\kappa + \sum_{i=1}^4 |k_i|^\kappa) dk_{1234} \\
& \lesssim \int \theta(2^{-q}k_{[1234]})^2 \psi_0(k_{[12]}, k_{[34]}) \left(\frac{\varepsilon^\kappa}{|k_{[34]}| |k_{[12]}|^{5-\kappa}} + \frac{\varepsilon^\kappa}{|k_{[34]}|^{1-\kappa} |k_{[12]}|^5} \right) dk_{[12][34]} \\
& \lesssim \int \theta(2^{-q}k)^2 2^{-q(2+\kappa)} \frac{\varepsilon^\kappa}{|k|^{1-2\kappa}} dk \lesssim \varepsilon^\kappa 2^{q\kappa},
\end{aligned}$$

where we used Lemma 6.1 in the second inequality and $|k_{[12]}| \gtrsim 2^q$ on the support of $\theta(2^{-q}k_{[1234]}) \psi_0(k_{[12]}, k_{[34]})$ in the third inequality. Now we have that for $\kappa > 0$ small enough

$$E[|\Delta_q(I_t^1 - \bar{I}_t^1)|^2] \lesssim 2^{q\kappa} \varepsilon^\kappa.$$

By a similar calculation as above and Gaussian hypercontractivity and Lemma 2.1 we obtain that for $\delta > 0$, $p > 1$

$$\pi_{0,\diamond}(K^\varepsilon, (u_1^\varepsilon)^{\diamond 2}) - \pi_{0,\diamond}(\bar{K}^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond 2}) \rightarrow 0 \text{ in } L^p(\Omega; C_T \mathcal{C}^{-\delta}).$$

Convergence for $\pi_0(u_2^\varepsilon, u_1^\varepsilon) - \pi_0(\bar{u}_2^\varepsilon, \bar{u}_1^\varepsilon)$

In this part we focus on $\pi_0(u_2^\varepsilon, u_1^\varepsilon)$ and prove that $\pi_0(u_2^\varepsilon, u_1^\varepsilon) - \pi_0(\bar{u}_2^\varepsilon, \bar{u}_1^\varepsilon) \rightarrow 0$ in $C_T \mathcal{C}^{-\delta}$. Now we have the following identity: for $t \in [0, T]$,

$$\pi_0(\bar{u}_2^\varepsilon, \bar{u}_1^\varepsilon)(t) - \pi_0(u_2^\varepsilon, u_1^\varepsilon)(t) = I_t^1 + 3I_t^2 - [\bar{I}_t^1 + 3\bar{I}_t^2] + J_t^1 + 3J_t^2.$$

Here

$$\begin{aligned}
I_t^1 &= 2^{-\frac{9}{2}} \int e_{k_{[1234]}} \psi_0(k_{[123]}, k_4) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[123]}) P_{\sigma-s_1}^\varepsilon(k_1) P_{\sigma-s_2}^\varepsilon(k_2) P_{\sigma-s_3}^\varepsilon(k_3) P_{t-s_4}^\varepsilon(k_4) W(d\eta_{1234}), \\
I_t^2 &= 2^{-\frac{9}{2}} \int \int e_{k_{[23]}} \psi_0(k_{[123]}, k_1) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[123]}) P_{\sigma-s_2}^\varepsilon(k_2) P_{\sigma-s_3}^\varepsilon(k_3) V_{t-\sigma}^\varepsilon(k_1) dk_1 W(d\eta_{23}),
\end{aligned}$$

and for $i = 1, 2$, \bar{I}_t^i is defined with $P_{t-\sigma}^\varepsilon(k_{[123]})$ replaced by $p_{t-\sigma}^\varepsilon(k_{[123]})$ and other $P^\varepsilon, V^\varepsilon$ replaced by $\bar{P}^\varepsilon, \bar{V}^\varepsilon$ respectively and for $i = 1, 2$, J_t^i is defined similar as I_t^i with each $k_{[123]}, e_{k_{[1234]}}, e_{k_{[23]}}$ replaced by $\tilde{k}_{[123]}, e_{\tilde{k}_{[1234]}}, e_{\tilde{k}_{[23]}}$.

Terms in the second chaos: First we consider I_t^2 and we have the following calculations:

$$\begin{aligned}
& E|\Delta_q(I_t^2 - \bar{I}_t^2)|^2 \\
& \lesssim \int \psi_0(k_{[123]}, k_1) \psi_0(k_{[234]}, k_4) \theta(2^{-q}k_{[23]})^2 \\
& \quad \frac{|\varepsilon k_{[123]}|^{\kappa/2} |\varepsilon k_{[234]}|^{\kappa/2} + |\varepsilon k_1|^{\kappa/2} |\varepsilon k_4|^{\kappa/2} + |\varepsilon k_2|^\kappa + |\varepsilon k_3|^\kappa}{|k_2|^2 |k_3|^2 |k_1|^2 (|k_1|^2 + |k_{[123]}|^2) |k_4|^2 (|k_4|^2 + |k_{[234]}|^2)} dk_{1234} \\
& \lesssim \varepsilon^\kappa \int 2^{-q(2-2\kappa)} \theta(2^{-q}k_{[23]})^2 \frac{1}{|k_2|^{2-\kappa} |k_3|^{2-\kappa}} dk_{23} \lesssim \varepsilon^\kappa 2^{3q\kappa},
\end{aligned}$$

where $\kappa > 0$ are small enough. Here we used (6.1, 6.2), Lemma 6.4 in the first inequality and $|k_{[123]}| \gtrsim 2^q$, $k_{[234]} \gtrsim 2^q$ in the second inequality and we used Lemma 6.1 in the last inequality. By a similar calculation as above, we know that

$$E|\Delta_q J_t^2|^2 \lesssim \int 2^{-q(2-2\kappa)} \theta(2^{-q} \tilde{k}_{[23]})^2 \frac{\varepsilon^\kappa}{|k_2|^{2-\kappa} |k_3|^2} dk_{23} \lesssim \varepsilon^\kappa 2^{3\kappa q}.$$

Here $\kappa > 0$ is small enough and in the first inequality we used $|k_{[123]}| \simeq N$ to deduce that $|k_i| \simeq N$ for some $i \in \{1, 2, 3\}$ and in the last inequality we used Lemmas 6.1 and 6.5.

Terms in the fourth chaos: Now for I_t^1 we have the following calculations:

$$\begin{aligned} & E[|\Delta_q(I_t^1 - \bar{I}_t^1)|^2] \\ & \lesssim \varepsilon^\kappa \int \frac{\theta(2^{-q} k_{[1234]})^2 \psi_0(k_{[123]}, k_4) (|k_{[123]}|^\kappa + \sum_{i=1}^4 |k_i|^\kappa)}{|k_1|^2 |k_2|^2 |k_3|^2 |k_4|^2 [|k_1|^2 + |k_2|^2 + |k_3|^2] |k_{[123]}|^2} dk_{1234} \\ & \lesssim \int 2^{-q(2-\kappa)} \theta(2^{-q} k)^2 \frac{\varepsilon^\kappa}{|k|} dk \lesssim \varepsilon^\kappa 2^{q\kappa}, \end{aligned}$$

where we used (6.1), (6.2) and Lemma 6.4 in the first inequality, Lemma 6.1 in the second inequality and $|k_{[123]}| \gtrsim 2^q$ in the third inequality. For J_t^1 , using Lemma 6.5 and by a similar argument we also obtain that

$$E|\Delta_q J_t^1|^2 \lesssim \varepsilon^\kappa 2^{q\kappa}.$$

Now by a similar calculation as above, Gaussian hypercontractivity and Lemma 2.1 we have that for $\delta > 0$, $p > 1$

$$\pi_{0,\diamond}(u_2^\varepsilon, u_1^\varepsilon) - \pi_{0,\diamond}(\bar{u}_2^\varepsilon, \bar{u}_1^\varepsilon) \rightarrow 0 \text{ in } L^p(\Omega; C_T \mathcal{C}^{-\delta}).$$

Convergence for $\pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond 2}) - \pi_{0,\diamond}(\bar{u}_2^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond 2})$ In this part we focus on $\pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond 2})$ and prove that $\pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond 2}) - \pi_{0,\diamond}(\bar{u}_2^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond 2}) \rightarrow 0$ in $C_T \mathcal{C}^{-1/2-\delta/2}$. Now we have the following identity: for $t \in [0, T]$,

$$\pi_{0,\diamond}(\bar{u}_2^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond 2}) - \pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond 2}) = I_t^1 + 6I_t^2 + 6I_t^3 - [\bar{I}_t^1 + 6\bar{I}_t^2 + 6\bar{I}_t^3] + J_t^1 + 6J_t^2 + 6J_t^3.$$

Here

$$\begin{aligned} I_t^1 &= 2^{-6} \int e_{k_{[12345]}} \psi_0(k_{[123]}, k_{[45]}) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[123]}) \Pi_{i=1}^3 P_{\sigma-s_i}^\varepsilon(k_i) \Pi_{i=4}^5 P_{t-s_i}^\varepsilon(k_i) W(d\eta_{12345}), \\ I_t^2 &= 2^{-6} \int e_{k_{[234]}} \psi_0(k_{[123]}, k_4 - k_1) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[123]}) \Pi_{i=2}^3 P_{\sigma-s_i}^\varepsilon(k_i) P_{t-s_4}^\varepsilon(k_4) V_{t-\sigma}^\varepsilon(k_1) W(d\eta_{234}), \\ I_t^3 &= 2^{-6} \int e_{k_3} \psi_0(k_{[123]}, k_{[12]}) \int_0^t d\sigma P_{\sigma-s_3}^\varepsilon(k_3) V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) P_{t-\sigma}^\varepsilon(k_{[123]}) W(d\eta_3), \end{aligned}$$

and for $i = 1, 2, 3$, \bar{I}_t^i is defined similarly with $P_{t-\sigma}^\varepsilon(k_{[123]})$ replaced by $p_{t-\sigma}(k_{[123]})$ and other $P^\varepsilon, V^\varepsilon$ replaced by $\bar{P}^\varepsilon, \bar{V}^\varepsilon$ respectively and for $i = 1, 2, 3$, J_t^i is defined similar as I_t^i with each $k_{[123]}, e_{k_{[12345]}}, e_{k_{[234]}}, e_{k_3}$ replaced by $\tilde{k}_{[123]}, e_{\tilde{k}_{[12345]}}, e_{\tilde{k}_{[234]}}, e_{\tilde{k}_3}$.

We consider the following term first:

$$I_t^3 - \bar{I}_t^3 - [\tilde{I}_t^3 - \tilde{\bar{I}}_t^3] + \tilde{I}_t^3 - \tilde{\bar{I}}_t^3 - C^\varepsilon(t) u_1^\varepsilon(t) - \bar{C}^\varepsilon(t) \bar{u}_1^\varepsilon(t),$$

where $\tilde{I}_t^3, \tilde{\tilde{I}}_t^3$ are defined similarly as I_t^3, \bar{I}_t^3 with $P_{\sigma-s_3}^\varepsilon(k_3), \bar{P}_{\sigma-s_3}^\varepsilon(k_3)$ replaced by $P_{t-s_3}^\varepsilon(k_3), \bar{P}_{t-s_3}^\varepsilon(k_3)$, respectively and $C(t)^\varepsilon = \frac{1}{2}[C_{11}^\varepsilon + \varphi_1^\varepsilon(t)], \bar{C}(t)^\varepsilon = \frac{1}{2}[\bar{C}_{11}^\varepsilon + \bar{\varphi}_1^\varepsilon(t)]$.

Since $\int |P_{t-s_3}^\varepsilon(k_3) - P_{\sigma-s_3}^\varepsilon(k_3)|^2 ds_3 \lesssim \frac{(t-\sigma)^{\kappa/2}}{|k_3|^{2-\kappa}}$ and

$$\int |P_{t-s_3}^\varepsilon(k_3) - P_{\sigma-s_3}^\varepsilon(k_3) - [\bar{P}_{t-s_3}^\varepsilon(k_3) - \bar{P}_{\sigma-s_3}^\varepsilon(k_3)]|^2 ds_3 \lesssim \frac{(t-\sigma)^{\kappa/2} \wedge \varepsilon^\kappa}{|k_3|^{2-\kappa}},$$

by a straightforward calculation we obtain that for $\kappa > 0$ small enough

$$\begin{aligned} & E[|\Delta_q(I_t^3 - \bar{I}_t^3 - [\tilde{I}_t^3 - \tilde{\tilde{I}}_t^3])|^2] \\ & \lesssim \int \theta(2^{-q}k_3)^2 \left[\frac{1}{|k_3|^{2-2\kappa}} \left(\int_0^t \int \varepsilon^{\kappa/2} (|k_{[123]}|^{\kappa/2} + |k_2|^{\kappa/2} + |k_1|^{\kappa/2}) \frac{e^{-(|k_{[123]}|^2 + |k_1|^2 + |k_2|^2)\bar{c}_f(t-\sigma)}}{|k_1|^2|k_2|^2} \right. \right. \\ & \quad \left. \left. (t-\sigma)^{\kappa/2} dk_{12} d\sigma \right)^2 + \frac{\varepsilon^\kappa}{|k_3|^{2-2\kappa}} \left(\int_0^t \int \frac{e^{-(|k_{[123]}|^2 + |k_1|^2 + |k_2|^2)(t-\sigma)}}{|k_1|^2|k_2|^2} (t-\sigma)^{\kappa/4} dk_{12} d\sigma \right)^2 \right] dk_3 \\ & \lesssim \varepsilon^\kappa 2^{q(1+3\kappa)}. \end{aligned}$$

Here in the last inequality we used that $\sup_{a \geq 0} a^r e^{-a} \leq C$ for $r \geq 0$ and Lemma 6.1. Moreover, by Lemmas 6.2 and 6.3 we obtain that

$$\begin{aligned} & E[|\Delta_q(\tilde{I}_t^3 - \tilde{\tilde{I}}_t^3 - u_1^\varepsilon(t)C^\varepsilon(t) + \bar{u}_1^\varepsilon(t)\bar{C}^\varepsilon(t))|^2] \\ & \lesssim \int \frac{1}{|k_3|^2} \theta(2^{-q}k_3) \left(\int \int_0^t |k_{[12]}|^{-\kappa} |k_3|^\kappa \right. \\ & \quad \left. (\varepsilon^{\kappa/2}|k_2|^{\kappa/2} + \varepsilon^{\kappa/2}|k_1|^{\kappa/2} + \varepsilon^{\kappa/2}|k_3|^{\kappa/2}) \frac{e^{-|k_1|^2(t-\sigma)\bar{c}_f - |k_2|^2(t-\sigma)\bar{c}_f}}{|k_1|^2|k_2|^2} dk_{12} d\sigma \right)^2 dk_3 \\ & \quad + \int \frac{\varepsilon^\kappa |k_3|^\kappa}{|k|^2} \theta(2^{-q}k_3)^2 \left(\int \int_0^t \frac{e^{-|k_2|^2(t-\sigma) - |k_1|^2(t-\sigma)}}{|k_1|^2|k_2|^2} |k_3|^\kappa |k_{[12]}|^{-\kappa} dk_{12} d\sigma \right)^2 dk_3 \\ & \lesssim \varepsilon^\kappa \int \theta(2^{-q}k_3) \frac{1}{|k_3|^{2-3\kappa}} dk_3 \lesssim \varepsilon^\kappa 2^{q(1+3\kappa)}. \end{aligned}$$

For J_t^3 we have

$$E[|\Delta_q J_t^3|^2] \lesssim \int \frac{1}{|k_3|^2} \theta(2^{-q}\tilde{k}_3) \left(\int \frac{1_{|k_1| \leq N, |k_2| \leq N}}{|k_1|^2|k_2|^2(|k_1|^2 + |k_2|^2 + |\tilde{k}_{[12]}|^2)} dk_{12} \right)^2 dk_3 \lesssim \varepsilon^\kappa 2^{q(1+3\kappa)}.$$

Here we used $2^q \simeq N$ in the last inequality.

Terms in the third chaos: Now we focus on the bounds for I_t^2 . We obtain the following inequalities:

$$\begin{aligned} & E|\Delta_q(I_t^2 - \bar{I}_t^2)|^2 \\ & \lesssim \int \theta(2^{-q}k_{[234]}) \psi_0(k_{[123]}, k_4 - k_1) \psi_0(k_{[235]}, k_4 - k_5) \\ & \quad \prod_{i=1}^5 \frac{1}{|k_i|^2} \frac{|k_{[123]}|^{\kappa/2} |k_{[235]}|^{\kappa/2} \varepsilon^\kappa + \sum_{i=2}^4 (\varepsilon |k_i|)^\kappa}{(|k_1|^2 + |k_{[123]}|^2 + |k_2|^2)(|k_5|^2 + |k_{[235]}|^2)} dk_{12345} \\ & \lesssim \int 2^{-q(1-\kappa)} \frac{\varepsilon^\kappa \theta(2^{-q}k_{[234]})}{|k_2|^{3-2\kappa} |k_3|^2 |k_4|^2} dk_{234} \lesssim \varepsilon^\kappa 2^{q(1+4\kappa)}, \end{aligned}$$

where we used Lemma 6.1 in the last inequality. For J_t^2 by a similar calculation as above and using the fact that $|k_{[235]}| \simeq N \gtrsim |\tilde{k}_{[235]}|$, we know that

$$E|\Delta_q J_t^2|^2 \lesssim \int 2^{-q(1-\kappa)} \theta(2^{-q} \tilde{k}_{[234]})^2 \frac{1}{|k_2|^3 |k_3|^2 |k_4|^2} dk_{234} \lesssim \varepsilon^\kappa 2^{(1+2\kappa)q}.$$

Here $\kappa > 0$ is small enough and in the last inequality we used Lemmas 6.1 and 6.6.

Term in the fifth chaos: Now we focus on the bounds for I_t^1 . We obtain the following inequalities:

$$\begin{aligned} & E|\Delta_q(I_t^1 - \tilde{I}_t^1)|^2 \\ & \lesssim \int \theta(2^{-q} k_{[12345]})^2 \psi_0(k_{[123]}, k_{[45]})^2 \Pi_{i=1}^5 \frac{1}{|k_i|^2} \frac{(\sum_{i=1}^5 |\varepsilon k_i|^\kappa + |\varepsilon k_{[123]}|^\kappa)}{|k_{[123]}|^2 (|k_1|^2 + |k_2|^2 + |k_{[123]}|^2)} dk_{12345} \\ & \lesssim \varepsilon^\kappa 2^{q(1+2\kappa)}. \end{aligned}$$

For J_t^1 by similar calculations for I_t^1 and using the fact that $|k_{[123]}| \simeq N \gtrsim |\tilde{k}_{[123]}|$ we also obtain that

$$E|\Delta_q J_t^1|^2 \lesssim \varepsilon^\kappa 2^{q(1+2\kappa)}.$$

By a similar calculation as above we also obtain that there exist $\kappa, \epsilon, \gamma > 0$ small enough such that

$$\begin{aligned} & E[|\Delta_q(\pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond 2})(t_1) - \pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond 2})(t_2) - \pi_{0,\diamond}(\bar{u}_2^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond 2})(t_1) + \pi_{0,\diamond}(\bar{u}_2^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond 2})(t_2))|^2] \\ & \lesssim \varepsilon^\gamma |t_1 - t_2|^\kappa 2^{q(1+\epsilon)}, \end{aligned}$$

which by Gaussian hypercontractivity and Lemma 2.1 implies that for every $\delta > 0, p > 1$, $\pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\diamond 2}) - \pi_{0,\diamond}(\bar{u}_2^\varepsilon, (\bar{u}_1^\varepsilon)^{\diamond 2}) \rightarrow 0$ in $L^p(\Omega; C_T \mathcal{C}^{-1/2-\delta/2})$.

6.2 Convergence of the error terms

In this subsection we prove $E_W^\varepsilon \rightarrow^P 0$ as $\varepsilon \rightarrow 0$.

Convergence for $\pi_0(K^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond 2})$

Now we have the following identity: for $t \in [0, T]$,

$$\pi_0(K^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond 2}) = I_t^1 + 4I_t^2 + 2I_t^3,$$

$$I_t^1 = 2^{-6} \int e_{\tilde{k}_{[1234]}} \psi_0(k_{[12]}, \tilde{k}_{[34]}) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[12]}) P_{\sigma-s_1}^\varepsilon(k_1) P_{\sigma-s_2}^\varepsilon(k_2) P_{t-s_3}^\varepsilon(k_3) P_{t-s_4}^\varepsilon(k_4) W(d\eta_{1234}),$$

$$I_t^2 = 2^{-6} \int \int e_{\tilde{k}_{[23]}} \psi_0(k_{[12]}, \tilde{k}_3 - k_1) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[12]}) P_{\sigma-s_2}^\varepsilon(k_2) P_{t-s_3}^\varepsilon(k_3) V_{t-\sigma}^\varepsilon(k_1) dk_1 W(d\eta_{23})$$

$$I_t^3 = 2^{-6} \int e_N^{i_1 i_2 i_3} \psi_0(k_{[12]}, -\tilde{k}_{[12]}) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[12]}) V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) dk_{12}.$$

Term in the 0-th chaos: We have

$$E[|\Delta_q I_t^3|^2] \lesssim \left(\int \frac{1_{|k_{[12]}| \simeq N \simeq 2^q} \psi_0(k_{[12]}, \tilde{k}_{[12]})}{|k_{[12]}|^3} dk_{[12]} \right)^2 \lesssim \varepsilon^\kappa 2^{q(3\kappa)}.$$

Term in the second chaos: Now we consider I_t^2 and we have the following calculations:

$$\begin{aligned} E|\Delta_q I_t^2|^2 &\lesssim \int \psi_0(k_{[12]}, \tilde{k}_3 - k_1) \psi_0(k_{[24]}, \tilde{k}_3 - k_4) \theta(2^{-q} \tilde{k}_{[23]})^2 \\ &\quad \frac{1}{|k_2|^2 |k_3|^2 |k_1|^2 (|k_1|^2 + |k_{12}|^2) |k_4|^2 (|k_4|^2 + |k_{[24]}|^2)} dk_{1234} \\ &\lesssim \int 2^{(-2+\kappa)q} \theta(2^{-q} \tilde{k}_{[23]})^2 \frac{1}{|k_2|^2 |k_3|^2} dk_{23} \lesssim \varepsilon^\kappa 2^{q2\kappa}, \end{aligned}$$

where $\kappa > 0$ is small enough. Here we used $|k_{[12]}| \gtrsim 2^q$ in the second inequality and used Lemmas 6.1, 6.6 in the third inequality.

Term in the fourth chaos: Now for I_t^1 we have the following calculations:

$$\begin{aligned} E[|\Delta_q I_t^1|^2] &\lesssim \int \theta(2^{-q} \tilde{k}_{[1234]})^2 \psi_0(k_{[12]}, \tilde{k}_{[34]}) \frac{1}{|k_{[34]}| |k_{[12]}|^{5-\kappa}} dk_{[12][34]} \\ &\lesssim 2^{-2q} \int \theta(2^{-q} \tilde{k}_{[1234]})^2 \frac{1}{|k_{[34]}| |k_{[12]}|^{3-\kappa}} dk_{[12][34]} \lesssim \varepsilon^\kappa 2^{q\kappa}, \end{aligned}$$

where we used Lemmas 6.1, 6.6 in the last inequality. By a similar calculation as above, Gaussian hypercontractivity and Lemma 2.1 we obtain that for $\delta > 0$ small enough, $p > 1$

$$\pi_0(K^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}) \rightarrow 0 \text{ in } L^p(\Omega; C_T \mathcal{C}^{-\delta}).$$

Convergence of $\pi_0(K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2})$ Now we have the following identity: for $t \in [0, T]$,

$$\pi_0(K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2}) = I_t^1 + 4I_t^2 + 2I_t^3,$$

$$\begin{aligned} I_t^1 &= 2^{-6} \int e_{\tilde{k}_{[1234]}} \psi_0(\tilde{k}_{[12]}, k_{[34]}) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(\tilde{k}_{[12]}) P_{\sigma-s_1}^\varepsilon(k_1) P_{\sigma-s_2}^\varepsilon(k_2) P_{t-s_3}^\varepsilon(k_3) P_{t-s_4}^\varepsilon(k_4) W(d\eta_{1234}), \\ I_t^2 &= 2^{-6} \int \int e_{\tilde{k}_{[23]}} \psi_0(\tilde{k}_{[12]}, k_3 - k_1) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(\tilde{k}_{[12]}) P_{\sigma-s_2}^\varepsilon(k_2) P_{t-s_3}^\varepsilon(k_3) V_{t-\sigma}^\varepsilon(k_1) dk_1 W(d\eta_{23}) \\ I_t^3 &= 2^{-6} \int e_N^{i_1 i_2 i_3} \psi_0(\tilde{k}_{[12]}, -k_{[12]}) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(\tilde{k}_{[12]}) V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) dk_{12}. \end{aligned}$$

I_t^3, I_t^2 can be estimated similarly as that for $\pi_0(K^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})$ and we only consider **Terms in the fourth chaos:** Now for I_t^1 we have the following calculations:

$$\begin{aligned} E|\Delta_q I_t^1|^2 &\lesssim \int \psi_0(\tilde{k}_{[12]}, k_{[34]}) \theta(2^{-q} \tilde{k}_{[1234]})^2 \frac{1}{|k_2|^2 |k_3|^2 |k_1|^2 (|k_1|^2 + |k_2|^2) |k_4|^2 |\tilde{k}_{[12]}|^2} dk_{1234} \\ &\lesssim \int 2^{-2q} \theta(2^{-q} \tilde{k}_{[1234]})^2 \frac{1}{|k_{[12]}|^{3-\kappa} |k_{[34]}|} dk_{[12][34]} \lesssim \varepsilon^\kappa 2^{q2\kappa}, \end{aligned}$$

where we used Lemmas 6.1, 6.6 in the last inequality. By a similar calculation as above, Gaussian hypercontractivity and Lemma 2.1 we obtain that for $\delta > 0$, $p > 1$

$$\pi_{0,\diamond}(K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2}) \rightarrow 0 \text{ in } L^p(\Omega; C_T \mathcal{C}^{-\delta}).$$

Convergence of $\pi_{0,\diamond}(K_1^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2})$

We have

$$\pi_0(K_1^\varepsilon, e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond, 2}) = I_t^1 + 4I_t^2 + 2I_t^3.$$

Here $I_t^i, i = 1, 2$ is defined similarly as that for $\pi_0(K^\varepsilon, e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond, 2})$ with $k_{[12]}, e_{\tilde{k}_{[1234]}}$ and $e_{\tilde{k}_{[23]}}$ replaced by $\tilde{k}_{[12]}, e_{\tilde{k}_{[1234]}}$ and $e_{\tilde{k}_{[23]}}$, respectively and

$$I_t^3 = 2^{-3} \int e_N^{i_1 i_2 i_3} e_N^{i'_1 i'_2 i'_3} \psi_0(\tilde{k}_{[12]}^{i'_1 i'_2 i'_3}, \widetilde{-k_{[12]}^{i_1 i_2 i_3}}) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(\tilde{k}_{[12]}^{i'_1 i'_2 i'_3}) V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) dk_{12},$$

for $i_j, i'_j \in \{-1, 0, 1\}$ for $j = 1, 2, 3$ with $\sum_j i_j^2 \neq 0, \sum_j (i'_j)^2 \neq 0$. Choose

$$C_{12}^{\varepsilon, i_1 i_2 i_3} = 2^{-5} \int_{-\infty}^t d\sigma P_{t-\sigma}^\varepsilon(\widetilde{-k_{[12]}^{i_1 i_2 i_3}}) V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) dk_{12},$$

and $\varphi_2^{\varepsilon, i_1 i_2 i_3}(t) = -2^{-5} \int_{-\infty}^0 d\sigma P_{t-\sigma}^\varepsilon(\widetilde{-k_{[12]}^{i_1 i_2 i_3}}) V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) dk_{12}$, we could easily obtain that

$$|C_{12}^{\varepsilon, i_1 i_2 i_3}| \simeq 1, \quad \sup_{t \in [0, T]} t^\rho |\varphi_2^{\varepsilon, i_1 i_2 i_3}(t)| \lesssim \varepsilon^\kappa,$$

for every $\rho > \kappa/2 > 0$. For the terms in $2I_t^3 - C_{12}^{\varepsilon, i_1 i_2 i_3} - \varphi_2^{\varepsilon, i_1 i_2 i_3}$, we know that $e_N^{i_1 i_2 i_3} e_N^{i'_1 i'_2 i'_3} \neq 1$ and we could easily obtain

$$E[|\Delta_q(2I_t^3 - C_{12}^\varepsilon - \varphi_2^\varepsilon)|^2] \lesssim \varepsilon^\kappa 2^{q(3\kappa)}.$$

Term in the second chaos: Now we consider I_t^2 and we have the following calculations:

$$\begin{aligned} E|\Delta_q I_t^2|^2 &\lesssim \int \psi_0(\tilde{k}_{[12]}, \tilde{k}_3 - k_1) \psi_0(\tilde{k}_{[24]}, \tilde{k}_3 - k_4) \theta(2^{-q} \tilde{k}_{[23]})^2 1_{|k_{[12]}| > N, |k_{[24]}| > N} \\ &\quad \frac{1}{|k_2|^2 |k_3|^2 |k_1|^2 (|k_1|^2 + |\tilde{k}_{12}|^2) |k_4|^2 (|k_4|^2 + |\tilde{k}_{[24]}|^2)} dk_{1234} \\ &\lesssim \varepsilon^\kappa \int 2^{(-2+2\kappa)q} \theta(2^{-q} \tilde{k}_{[23]})^2 \frac{1}{|k_2|^{2-\kappa} |k_3|^2} dk_{23} \lesssim \varepsilon^\kappa 2^{3q\kappa}, \end{aligned}$$

where $\kappa > 0$ is small enough. Here we used $|k_i| \simeq N$ for some $i \in \{1, 2, 4\}$ in the second inequality and Lemmas 6.1, 6.6 in the third inequality.

Term in the fourth chaos: Now for I_t^1 we have the following calculations:

$$\begin{aligned} E[|\Delta_q I_t^1|^2] &\lesssim \varepsilon^\kappa \int \theta(2^{-q} \tilde{k}_{[1234]})^2 \psi_0(\tilde{k}_{[12]}, \tilde{k}_{[34]}) \frac{1}{|k_{[34]}| |\tilde{k}_{[12]}|^{5-\kappa}} dk_{[12][34]} \\ &\lesssim \varepsilon^\kappa 2^{-2q} \int \theta(2^{-q} \tilde{k}_{[1234]})^2 \frac{1}{|k_{[34]}| |\tilde{k}_{[12]}|^{3-\kappa}} dk_{[12][34]} \lesssim \varepsilon^\kappa 2^{q\kappa}, \end{aligned}$$

where we used Lemmas 6.1, 6.5 in the last inequality. By a similar calculation as above, Gaussian hypercontractivity and Lemma 2.1 we obtain that for $\delta > 0$ small enough, $p > 1$

$$\pi_{0, \diamond}(K_1^\varepsilon, e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond, 2}) \rightarrow 0 \text{ in } L^p(\Omega; C_T \mathcal{C}^{-\delta}).$$

Convergence of $(u_1^\varepsilon)^{\diamond,2} e_N^{i_1 i_2 i_3}$ By a similar calculation as that for $(u_1^\varepsilon)^{\diamond,2}$ we know that

$$E|\Delta_q[(u_1^\varepsilon)^{\diamond,2} e_N^{i_1 i_2 i_3}]|^2 \lesssim \int \theta(2^{-q} \tilde{k}_{[12]})^2 \frac{1}{|k_1|^2 |k_2|^2} dk_1 dk_2 \lesssim \varepsilon^\kappa 2^{(\kappa+2)q}.$$

Here $\kappa > 0$ is small enough and in the last inequality we used Lemmas 6.1, 6.6. Then by Gaussian hypercontractivity and Lemma 2.1 we obtain that for every $\delta > 0, p > 1$, $(u_1^\varepsilon)^{\diamond,2} e_N^{i_1 i_2 i_3} \rightarrow 0$ in $L^p(\Omega; C_T \mathcal{C}^{-1-\delta})$.

Convergence of $\pi_0(u_2^\varepsilon, e_N^{i_1 i_2 i_3} u_1^\varepsilon)$ Now we have the following identity: for $t \in [0, T]$,

$$\pi_0(u_2^\varepsilon, e_N^{i_1 i_2 i_3} u_1^\varepsilon)(t) = -[I_t^1 + 3I_t^2 + J_t^1 + 3J_t^2].$$

Here

$$\begin{aligned} I_t^1 &= 2^{-\frac{9}{2}} \int e_{\tilde{k}_{[1234]}} \psi_0(k_{[123]}, \tilde{k}_4) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[123]}) P_{\sigma-s_1}^\varepsilon(k_1) P_{\sigma-s_2}^\varepsilon(k_2) P_{\sigma-s_3}^\varepsilon(k_3) P_{t-s_4}^\varepsilon(k_4) W(d\eta_{1234}), \\ I_t^2 &= 2^{-\frac{9}{2}} \int \int e_{\tilde{k}_{[23]}} \psi_0(k_{[123]}, \tilde{k}_1) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[123]}) P_{\sigma-s_2}^\varepsilon(k_2) P_{\sigma-s_3}^\varepsilon(k_3) V_{t-\sigma}^\varepsilon(k_1) dk_1 W(d\eta_{23}), \end{aligned}$$

and for $i = 1, 2$, J_t^i is defined similarly as I_t^3 with each $k_{[123]}, e_{\tilde{k}_{[1234]}}, e_{\tilde{k}_{[23]}}$ replaced by $\tilde{k}_{[123]}, e_{\tilde{k}_{[1234]}}, e_{\tilde{k}_{[23]}}$.

Term in the second chaos: First we consider I_t^2 and by a similar calculations as that for $\pi_0(u_2^\varepsilon, u_1^\varepsilon)$:

$$E|\Delta_q I_t^2|^2 \lesssim \varepsilon^\kappa \int 2^{-q(2-2\kappa)} \theta(2^{-q} \tilde{k}_{[23]})^2 \frac{1}{|k_2|^2 |k_3|^2} dk_{23} \lesssim \varepsilon^\kappa 2^{q4\kappa},$$

where $\kappa > 0$ are small enough and we used $|k_{[123]}| \simeq |\tilde{k}_1| \simeq N$ in the first inequality, we used Lemmas 6.1, 6.5 in the last inequality. By a similar calculation as above, we know that

$$E|\Delta_q J_t^2|^2 \lesssim \varepsilon^\kappa \int 2^{-q(2-2\kappa)} \theta(2^{-q} \tilde{k}_{[23]})^2 \frac{1}{|k_2|^2 |k_3|^2} dk_{23} \lesssim \varepsilon^\kappa 2^{3\kappa q}.$$

Here $\kappa > 0$ is small enough and we used $|\tilde{k}_{[123]}| \simeq |\tilde{k}_1| \simeq N$ in the first inequality, we used Lemmas 6.1, 6.5 in the last inequality.

Terms in the fourth chaos: Now for I_t^1, J_t^1 we have similar estimates:

$$E[|\Delta_q I_t^1|^2 + |\Delta_q J_t^1|^2] \lesssim \varepsilon^\kappa 2^{\kappa q}.$$

Here for I_t^1 we used $|k_{[123]}| \simeq \tilde{k}_4 \simeq N$ and for J_t^1 we used $|k_{[123]}| \simeq N \gtrsim \tilde{k}_{[123]}$. Now by a similar calculation as above, Gaussian hypercontractivity and Lemma 2.1 we have that for $\delta > 0, p > 1$

$$\pi_0(u_2^\varepsilon, e_N^{i_1 i_2 i_3} u_1^\varepsilon) \rightarrow 0 \text{ in } L^p(\Omega; C_T \mathcal{C}^{-\delta}).$$

Convergence of $\pi_{0,\diamond}(u_2^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2})$ Now we have the following identity: for $t \in [0, T]$,

$$\pi_{0,\diamond}(u_2^\varepsilon, e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond,2}) = -[I_t^1 + 6I_t^2 + 6I_t^3 + J_t^1 + 6J_t^2 + 6J_t^3]$$

$$\begin{aligned}
I_t^1 &= 2^{-6} \int e_{\tilde{k}_{[12345]}} \psi_0(k_{[123]}, \tilde{k}_{[45]}) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[123]}) \Pi_{i=1}^3 P_{\sigma-s_i}^\varepsilon(k_i) \Pi_{i=4}^5 P_{t-s_i}^\varepsilon(k_i) W(d\eta_{12345}), \\
I_t^2 &= 2^{-6} \int \int e_{\tilde{k}_{[234]}} \psi_0(k_{[123]}, \tilde{k}_4 - k_1) \int_0^t d\sigma P_{t-\sigma}^\varepsilon(k_{[123]}) \Pi_{i=2}^3 P_{\sigma-s_i}^\varepsilon(k_i) P_{t-s_4}^\varepsilon(k_4) V_{t-\sigma}^\varepsilon(k_1) dk_1 W(d\eta_{234}), \\
I_t^3 &= 2^{-6} \int \int e_{\tilde{k}_3} \psi_0(k_{[123]}, \tilde{k}_{[12]}) \int_0^t d\sigma P_{\sigma-s_3}^\varepsilon(k_3) V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) P_{t-\sigma}^\varepsilon(k_{[123]}) dk_{12} W(d\eta_3),
\end{aligned}$$

and for $i = 1, 2, 3$, J_t^i is defined similarly as I_t^i with each $k_{[123]}, e_{\tilde{k}_{[12345]}}, e_{\tilde{k}_{[234]}}, e_{\tilde{k}_3}$ replaced by $\tilde{k}_{[123]}, e_{\tilde{k}_{[12345]}}, e_{\tilde{k}_{[234]}}, e_{\tilde{k}_3}$.

Terms in the first chaos We consider J_t^3 and I_t^3 can be estimated similarly: $J_t^3 = J_t^{31} + J_t^{32}$, with J_t^{31}, J_t^{32} associated with the terms that $\tilde{k}_3 \neq k_3$ and $\tilde{k}_3 = k_3$, respectively. For J_t^{31} we have

$$\begin{aligned}
&E[|\Delta_q J_t^{31}|^2] \\
&\lesssim \int \frac{1_{|k_3| \lesssim N}}{|k_3|^2} \theta(2^{-q} \tilde{k}_3) \left(\int \frac{1_{|k_1| \lesssim N, |k_2| \lesssim N}}{|k_1|^2 |k_2|^2 (|k_1|^2 + |k_2|^2 + |\tilde{k}_{[123]}|^2)} dk_{12} \right)^2 dk_3 \lesssim \varepsilon^\kappa 2^{q(1+3\kappa)}.
\end{aligned}$$

Here we used $|\tilde{k}_3| \simeq 2^q \simeq N$ in the last inequality. For J_t^{32} we have

$$J_t^{32} - \tilde{J}_t^{32} + \tilde{J}_t^{32} - C_2^\varepsilon(t) u_1^\varepsilon(t),$$

with \tilde{J}_t^{32} defined as J_t^{32} with $P_{\sigma-s_3}^\varepsilon(k_3)$ replaced by $P_{t-s_3}^\varepsilon(k_3)$ and with $C_2^\varepsilon(t) = \frac{1}{2}(C_{12}^\varepsilon + \varphi_2^\varepsilon(t))$.

Since $\int |P_{t-s_3}^\varepsilon(k_3) - P_{\sigma-s_3}^\varepsilon(k_3)|^2 ds_3 \leq C \frac{(t-\sigma)^{\kappa/2}}{|k_3|^{2-\kappa}}$, by a straightforward calculation we obtain that for $\kappa > 0$ small enough

$$\begin{aligned}
&E[|\Delta_q(J_t^{32} - \tilde{J}_t^{32})|^2] \\
&\lesssim \int \theta(2^{-q} k_3)^2 \frac{1}{|k_3|^{2-4\kappa}} \left(\int_0^t \int \frac{e^{-(|\tilde{k}_{[123]}|^2 + |k_1|^2 + |k_2|^2) \bar{c}_f(t-\sigma)}}{|k_1|^2 |k_2|^2} \right. \\
&\quad \left. (t-\sigma)^\kappa dk_{12} d\sigma \right)^2 dk_3 \lesssim \varepsilon^\kappa 2^{q(1+5\kappa)}.
\end{aligned}$$

Here in the last inequality we used that $|k_{123}| \simeq N$ implies that $|k_i| \simeq N$ for some $i \in \{1, 2, 3\}$ and $\sup_{a \geq 0} a^r e^{-a} \leq C$ for $r \geq 0$ and Lemma 6.1. Moreover, by Lemmas 6.2 and 6.3 we obtain that

$$\begin{aligned}
&E[|\Delta_q(\tilde{J}_t^{32} - u_1^\varepsilon(t) C_2^\varepsilon(t))|^2] \\
&\lesssim \int \frac{1}{|k_3|^2} \theta(2^{-q} k_3) \left(\int \int_0^t |\tilde{k}_{[12]}|^{-\kappa} |k_3|^\kappa \frac{e^{-|k_1|^2(t-\sigma)\bar{c}_f - |k_2|^2(t-\sigma)\bar{c}_f}}{|k_1|^2 |k_2|^2} dk_{12} d\sigma \right)^2 dk_3 \\
&\lesssim \int \theta(2^{-q} k_3) \frac{1}{|k_3|^{2-2\kappa}} dk_3 \left[\int_{|k_{[12]}| \leq N} \frac{1}{|\tilde{k}_{[12]}|^\kappa |k_{[12]}|^3} dk_{[12]} + \varepsilon^{\kappa/2} \int_{|k_{[12]}| > N} \frac{1}{|\tilde{k}_{[12]}|^{3+\kappa/2}} dk_{[12]} \right]^2 \lesssim \varepsilon^\kappa 2^{q(1+3\kappa)},
\end{aligned}$$

where in the last inequality we used that if $|k_{[12]}| \leq N$, $|\tilde{k}_{[12]}| \simeq N$.

Terms in the third and fifth chaos can be estimated similarly as that for $\pi_{0,\diamond}(u_2^\varepsilon, (u_1^\varepsilon)^{\odot 2})$ and we also obtain that there exist $\kappa, \epsilon, \gamma > 0$ small enough such that

$$\begin{aligned}
&E[|\Delta_q(\pi_{0,\diamond}(u_2^\varepsilon, e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\odot 2})(t_1) - \pi_{0,\diamond}(u_2^\varepsilon, e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\odot 2})(t_2))|^2] \\
&\lesssim \varepsilon^\gamma |t_1 - t_2|^\kappa 2^{q(1+\epsilon)},
\end{aligned}$$

which by Gaussian hypercontractivity and Lemma 2.1 implies that for every $\delta > 0, p > 1$, $\pi_{0,\diamond}(u_2^\varepsilon, e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond,2}) \rightarrow 0$ in $L^p(\Omega; C_T \mathcal{C}^{-1/2-\delta/2})$.

6.3 Convergence of the random operator

The purpose of this subsection is to prove that A_N defined in Proposition 4.3 converges to zero in probability. Here we follow essentially the same arguments as [GP15, Section 10.2].

Theorem 6.7 For every $T \geq 0, 0 < \eta < \delta, r \geq 1$ we have

$$E[(A_N)^r]^{1/r} \lesssim N^{-\delta+\eta}.$$

By a similar argument as [GP15, Section 10.2] we have the following lemmas.

Lemma 6.8 We have

$$\begin{aligned} & A_N^1(K^\varepsilon + K_1^\varepsilon, (u_1^\varepsilon)^{\diamond,2} + e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond,2})(f) + A_N^2(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon, (u_1^\varepsilon)^{\diamond,2} + e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond,2})(f) \\ &= \sum_{p,q} \int_{\mathbb{T}^3} g_{p,q}^N(t, x, y) \Delta_p f(y) dy \end{aligned}$$

with

$$\mathcal{F}g_{p,q}^N(t, x, \cdot)(k) = \sum_{k_1, k_2} \Gamma_{p,q}^N(x, k, k_1, k_2) \mathcal{F}(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon)(t, k_1) \mathcal{F}((u_1^\varepsilon)^{\diamond,2} + e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond,2})(t, k_2).$$

Here

$$\begin{aligned} \Gamma_{p,q}^N(x, k, k_1, k_2) &= 2^{-9/2} e^{i(k_1+k_2-k)\pi x} \theta_q(k_1 + k_2 - k) \tilde{\theta}_p(k) \psi_<(k, k_1) \psi_0(k_1 - k, k_2) \\ &\quad (1_{|k_1-k|_\infty > N} 1_{|k_1|_\infty \leq N} + 1_{|k_1-k|_\infty \leq N} 1_{N < |k_1|_\infty \leq 3N}), \end{aligned}$$

with $\tilde{\theta}_p$ be a smooth function supported in an annulus $2^p \mathcal{A}$ such that $\tilde{\theta}_p \theta_p = \theta_p$.

Lemma 6.9 For all $r \geq 1$ we have

$$\begin{aligned} & E[|(A_N^1(K^\varepsilon + K_1^\varepsilon, (u_1^\varepsilon)^{\diamond,2} + e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond,2})) + A_N^2(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon, (u_1^\varepsilon)^{\diamond,2} + e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond,2}))(t) \\ & - (A_N^1(K^\varepsilon + K_1^\varepsilon, (u_1^\varepsilon)^{\diamond,2} + e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond,2})) + A_N^2(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon, (u_1^\varepsilon)^{\diamond,2} + e_N^{i_1 i_2 i_3}(u_1^\varepsilon)^{\diamond,2}))(s)|^r]_{L(C^{1-\delta}, B_{r,r}^{-1/2-2\delta+\kappa})} \\ & \lesssim \sum_{p,q} 2^{qr(-1/2-2\delta+\kappa)} 2^{-pr(1-\delta)} \left(\sup_{x \in \mathbb{T}^3} \sum_k E[|(\mathcal{F}g_{p,q}^N(t, x, \cdot) - \mathcal{F}g_{p,q}^N(s, x, \cdot))(k)|^2] \right)^{r/2}. \end{aligned}$$

Lemma 6.10 For all $p, q \geq -1$, all $0 \leq t_1 < t_2$, and all $\lambda, \kappa \in (0, 1]$ we have

$$\sum_k E[|(\mathcal{F}g_{p,q}^N(t_2, x, \cdot) - \mathcal{F}g_{p,q}^N(t_1, x, \cdot))(k)|^2] \lesssim 1_{2^p, 2^q \lesssim N} (2^{3p} 2^{2q} + 2^{2p} 2^{3q}) N^{-2+2\lambda+\kappa} |t_1 - t_2|^\lambda.$$

Proof We only prove the estimate for $\sum_k E[|\mathcal{F}g_{p,q}^N(t, x, \cdot)(k)|^2]$ and the result can be obtained by essentially the same arguments. First we consider $\mathcal{F}\tilde{K}^\varepsilon(t, l_1) \mathcal{F}(u_1^\varepsilon)^{\diamond,2}(t, l_2)$. We have the following chaos decomposition:

$$\mathcal{F}\tilde{K}^\varepsilon(t, l_1) \mathcal{F}(u_1^\varepsilon)^{\diamond,2}(t, l_2) = I_t^1 + I_t^2 + I_t^3.$$

Here

$$\begin{aligned}
I_t^1 &= 2^{-3} \int 1_{k_{[12]}=l_1, k_{[34]}=l_2} \int_0^t d\sigma p_{t-\sigma}^\varepsilon(k_{[12]}) \varphi(\varepsilon k_{[12]}) P_{\sigma-s_1}^\varepsilon(k_1) P_{\sigma-s_2}^\varepsilon(k_2) P_{t-s_3}^\varepsilon(k_3) P_{t-s_4}^\varepsilon(k_4) W(d\eta_{1234}), \\
I_t^2 &= 2^{-3} \int 1_{k_{[12]}=l_1, k_3-k_1=l_2} \int_0^t d\sigma p_{t-\sigma}^\varepsilon(k_{[12]}) \varphi(\varepsilon k_{[12]}) P_{\sigma-s_2}^\varepsilon(k_2) P_{t-s_3}^\varepsilon(k_3) V_{t-\sigma}^\varepsilon(k_1) dk_1 W(d\eta_{23}), \\
I_t^3 &= 2^{-3} \int \int_0^t d\sigma 1_{k_{[12]}=l_1, -k_{[12]}=l_2} V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) p_{t-\sigma}^\varepsilon(k_{[12]}) \varphi(\varepsilon k_{[12]}) dk_{12},
\end{aligned}$$

Term in the chaos of order 0 By a similar calculation as in Section 6.1 we have

$$\begin{aligned}
& \sum_k \left| \sum_{k_1, k_2} \Gamma_{p,q}^N(x, k, k_1, k_2) 1_{k_1+k_2=0} I_t^3 \right|^2 \\
& \lesssim \sum_k \left| \sum_{k_1} \Gamma_{p,q}^N(x, k, k_1, -k_1) \frac{1}{|k_1|^3} \right|^2 \\
& \lesssim \sum_k \tilde{\theta}_p(k)^2 \theta_q(-k)^2 \left| \sum_{k_1} (1_{|k_1-k|_\infty > N} 1_{|k_1|_\infty \leq N} + 1_{|k_1-k|_\infty \leq N} 1_{N < |k_1|_\infty \leq 3N}) \psi_{<}(k, k_1) \psi_0(k_1 - k, k_1) \frac{1}{|k_1|^3} \right|^2.
\end{aligned}$$

In the first case without loss of generality we assume that $|k_1^i - k^i| > N$ for some i . Then there are at most $|k^i|$ values of k_1^i with $|k_1^i| < N$ and $|k_1^i - k^i| > N$. In the second case without loss of generality we assume that $|k_1^i| > N$ for some i . Then there are at most $|k^i|$ values of k_1^i with $|k_1^i| > N$ and $|k_1^i - k^i| \leq N$. Moreover observe that $|k_1| \simeq N$ on the support of $(1_{|k_1-k|_\infty > N} 1_{|k_1|_\infty \leq N} + 1_{|k_1-k|_\infty \leq N} 1_{|k_1|_\infty > N}) \psi_0(k - k_1, k_1)$ and that $|k| \lesssim N$ whenever $1_{|k_1|_\infty \leq 3N} \psi_{<}(k, k_1) \neq 0$, which implies that the above term is bounded by

$$\sum_k \tilde{\theta}_p(k)^2 \theta_q(-k)^2 |k|^2 1_{|k| \lesssim N} N^{-2} \lesssim 1_{2^p, 2^q \lesssim N} 2^{3p} 2^{2q} N^{-2}.$$

Term in the second chaos By a similar calculation as in Section 6.1 we have

$$\begin{aligned}
& \sum_k E \left| \sum_{l_1, l_2} \Gamma_{p,q}^N(x, k, l_1, l_2) I_t^2 \right|^2 \\
& \lesssim \sum_k 1_{2^p, 2^q \lesssim N} \tilde{\theta}_p(k)^2 \int \theta_q(k_{[23]} - k)^2 \Pi_{i=2}^3 \frac{1}{|k_i|^2} \left[\int \psi_{<}(k, k_{[12]}) \frac{1}{(|k_{[12]}|^2 + |k_1|^2) |k_1|^2} \right. \\
& \quad \left. (1_{|k_{[12]}-k|_\infty > N, |k_{[12]}|_\infty \leq N} + 1_{|k_{[12]}-k|_\infty \leq N, N < |k_{[12]}|_\infty \leq 3N}) dk_1 \right]^2 dk_3 \\
& \lesssim \sum_k 1_{2^p, 2^q \lesssim N} \tilde{\theta}_p(k)^2 \int \theta_q(k_{[23]} - k)^2 \frac{1}{|k_{[23]}|} N^{-2+\kappa} dk_{[23]} \\
& \lesssim \sum_k 1_{2^p, 2^q \lesssim N} \tilde{\theta}_p(k)^2 N^{-2+\kappa} 2^{2q} \lesssim 1_{2^p, 2^q \lesssim N} 2^{3p} 2^{2q} N^{-2+\kappa}.
\end{aligned}$$

Here in the second inequality we used that $|k_{[12]}| \simeq N$ on the support of $1_{|k_{[12]}-k|_\infty > N, |k_{[12]}|_\infty \leq N} \psi_{<}(k, k_{[12]})$ and in the third inequality we used Lemma 6.5.

Term in the forth chaos We have

$$\begin{aligned}
& \sum_k E \left| \sum_{l_1, l_2} \Gamma_{p,q}^N(x, k, l_1, l_2) I_t^1 \right|^2 \\
& \lesssim \sum_k \tilde{\theta}_p(k)^2 \int \theta_q(k_{[1234]} - k)^2 \psi_{<}(k, k_{[12]})^2 \psi_0(k_{[12]} - k, k_{[34]})^2 \frac{1_{2^p, 2^q \lesssim N}}{|k_2|^2 |k_3|^2 |k_1|^2 |k_4|^2 |k_{[12]}|^4} \\
& \quad (1_{|k_{[12]} - k|_\infty > N, |k_{[12]}|_\infty \leq N} + 1_{|k_{[12]} - k|_\infty \leq N, N < |k_{[12]}|_\infty \leq 3N}) dk_{1234} \\
& \lesssim \sum_k \tilde{\theta}_p(k)^2 \int \theta_q(k_{[1234]} - k)^2 \frac{1_{2^p, 2^q \lesssim N}}{|k_{[1234]}|^{1-\kappa}} dk_{[1234]} N^{-2-\kappa} \\
& \lesssim 1_{2^p, 2^q \lesssim N} \sum_k \tilde{\theta}_p(k)^2 N^{-2-\kappa} 2^{(2+\kappa)q} \\
& \lesssim 1_{2^p, 2^q \lesssim N} 2^{3p} 2^{2q} N^{-2},
\end{aligned}$$

where we used Lemma 6.1 and $|k_{[12]}| \simeq N$ in the second inequality.

Moreover we consider

$$\mathcal{F} \tilde{K}_1^\varepsilon(t, k_1) \mathcal{F}(u_1^\varepsilon)^{\diamond, 2}(t, k_2) = J_t^1 + J_t^2 + J_t^3.$$

Here $J_t^i, i = 1, 2$, is defined similar as $I_t^i, i = 1, 2$, with $k_{[12]}, k_{[1234]}, k_{[23]}, e_{k_{[23]}}, e_{k_{1234}}$ replaced by $\tilde{k}_{[12]}, \tilde{k}_{[1234]}, \tilde{k}_{[23]}, e_{\tilde{k}_{1234}}$, respectively and

$$J_t^3 = 2^{-3} \int \int_0^t d\sigma 1_{\tilde{k}_{[12]}=l_1, -k_{[12]}=l_2} V_{t-\sigma}^\varepsilon(k_1) V_{t-\sigma}^\varepsilon(k_2) p_{t-\sigma}^\varepsilon(\tilde{k}_{[12]}) \varphi(\varepsilon \tilde{k}_{[12]}) dk_{12}.$$

Terms in the chaos of order 0 We have

$$\begin{aligned}
& \sum_k \left| \sum_{k_1, k_2} \Gamma_{p,q}^N(x, k, \tilde{k}_1, k_2) 1_{k_1+k_2=0} J_t^3 \right|^2 \\
& \lesssim \sum_k \left| \sum_{k_1} \Gamma_{p,q}^N(x, k, \tilde{k}_1, -k_1) 1_{|k_1|_\infty \lesssim N} \frac{1}{|k_1| |\tilde{k}_1|^2} \right|^2 \\
& \lesssim \sum_k \tilde{\theta}_p(k)^2 \theta_q(\tilde{k})^2 \left| \sum_{k_1} (1_{|\tilde{k}_1 - k|_\infty > N} 1_{|\tilde{k}_1|_\infty \leq N} + 1_{|\tilde{k}_1 - k|_\infty \leq N} 1_{N < |\tilde{k}_1|_\infty \leq 3N}) \psi_{<}(k, \tilde{k}_1) 1_{|k_1|_\infty \lesssim N} \frac{1}{|k_1| |\tilde{k}_1|^2} \right|^2.
\end{aligned}$$

Similarly as above we obtain there are at most $|k^i|$ values of \tilde{k}_1^i with $|\tilde{k}_1^i| > N$ and $|\tilde{k}_1^i - k^i| \leq N$ or $|\tilde{k}_1^i| > N$ and $|\tilde{k}_1^i - k^i| \leq N$. Moreover observe that $|\tilde{k}_1| \simeq N$ on the support of $1_{|\tilde{k}_1 - k|_\infty > N} 1_{|\tilde{k}_1|_\infty \leq N} \psi_{<}(k, \tilde{k}_1)$ and that $|k| \lesssim N$ whenever $1_{|k_1|_\infty < 3N} \psi_{<}(k, k_1) \neq 0$, which implies that the above term is bounded by

$$\sum_k \tilde{\theta}_p(k)^2 \theta_q(\tilde{k})^2 |k|^2 1_{|k| \lesssim N} N^{-2} \lesssim 1_{2^p, 2^q \lesssim N} 2^{2p} 2^{3q} N^{-2}.$$

For the terms in the second chaos by a similar calculation as above we obtain the estimates.

Terms in the forth chaos We have

$$\begin{aligned}
& \sum_k E \left| \sum_{l_1, l_2} \Gamma_{p,q}^N(x, k, l_1, l_2) J_t^3 \right|^2 \\
& \lesssim \sum_k \tilde{\theta}_p(k)^2 \int \theta_q(\tilde{k}_{[1234]} - k)^2 \psi_{<}(k, \tilde{k}_{[12]})^2 \psi_0(\tilde{k}_{[12]}, k_{[34]})^2 \frac{1_{2^p, 2^q \lesssim N}}{|k_2|^2 |k_3|^2 |k_1|^2 |k_4|^2 |\tilde{k}_{[12]}|^4} 1_{|k_{[12]}| \lesssim N, |k_{[34]}| \lesssim N} dk_{1234} \\
& \quad (1_{|\tilde{k}_{12}-k|_\infty > N, |\tilde{k}_{12}|_\infty \leq N} + 1_{|\tilde{k}_{12}-k|_\infty \leq N, N < |\tilde{k}_{12}|_\infty \leq 3N}) \\
& \lesssim 1_{2^p, 2^q \lesssim N} \sum_k \tilde{\theta}_p(k)^2 \int \theta_q(\tilde{k}_{[1234]} - k)^2 \frac{1}{|k_1|^2 |k_2|^2 |k_3|^2 |k_4|^2} N^{-4} 1_{|k_{[12]}| \lesssim N, |k_{[34]}| \lesssim N} dk_{1234} \\
& \lesssim 1_{2^p, 2^q \lesssim N} \sum_k \tilde{\theta}_p(k)^2 \int \theta_q(\tilde{k}_{[1234]} - k) \frac{1}{|k_{[12]}|^2 |k_{[34]}|^2} dk_{[12][34]} N^{-2} \\
& \lesssim 1_{2^p, 2^q \lesssim N} 2^{3p} 2^{2q} N^{-2}.
\end{aligned}$$

Here in the second inequality we used $|\tilde{k}_{[12]}| \simeq N$ and in the third inequality we used $N^{-1} \lesssim |k_{[12]}|^{-1}, N^{-1} \lesssim |k_{[34]}|^{-1}$ and in the last inequality we used Lemma 6.5.

Furthermore for the terms associated with $\tilde{K}^\varepsilon(u_1^\varepsilon)^{\diamond, 2} e_N^{i_1 i_2 i_3}$ and $\tilde{K}_1^\varepsilon(u_1^\varepsilon)^{\diamond, 2} e_N^{i_1 i_2 i_3}$ we can also obtain similar estimates by similar arguments. Thus the result follows. \square

Proof of Theorem 6.10 For $t, s \geq 0$ we have

$$\begin{aligned}
& E[|(A_N(K^\varepsilon + K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}) + B_N(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}))(t) \\
& \quad - (A_N(K^\varepsilon + K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}) + B_N(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}))(s)|^r]_{L(C^{1-\delta}, C^{-1/2-2\delta})}^{1/r} \\
& \lesssim E[|(A_N(K^\varepsilon + K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}) + B_N(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}))(t) \\
& \quad - (A_N(K^\varepsilon + K_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}) + B_N(\tilde{K}^\varepsilon + \tilde{K}_1^\varepsilon, (u_1^\varepsilon)^{\diamond, 2} + e_N^{i_1 i_2 i_3} (u_1^\varepsilon)^{\diamond, 2}))(s)|^r]_{L(C^{1-\delta}, B_{r,r}^{-1/2-2\delta+\kappa})}^{1/r} \\
& \lesssim \left[\sum_{p,q} 2^{qr(-1/2-2\delta+\kappa)} 2^{-pr(1-\delta)} 1_{2^p, 2^q \lesssim N} [(2^{3p} 2^{2q} + 2^{2p} 2^{3q}) |t - s|^\lambda N^{-2+2\lambda+\kappa}]^{r/2} \right]^{1/r} \\
& \lesssim |t - s|^{\lambda/2} N^{-\delta+2\kappa+\lambda}.
\end{aligned}$$

Here $\delta > 2\kappa + \lambda > 0$. Thus the result follows by using Kolmogorov's continuity criterion. \square

6.4 Convergence of D^N

In this subsection we prove that $D^N \xrightarrow{P} 0$ as $\varepsilon \rightarrow 0$. Now we have the following identity: for $t \in [0, T]$,

$$\pi_0((I - P_N)\pi_{<}(u_2^\varepsilon, K^\varepsilon), (u_1^\varepsilon)^{\diamond, 2})(t) + \pi_0(P_N\pi_{<}(u_2^\varepsilon, (P_{3N} - P_N)\tilde{K}^\varepsilon), (u_1^\varepsilon)^{\diamond, 2})(t) = \sum_{i=1}^4 (I_t^i + J_t^i).$$

Here

$$\begin{aligned}
I_t^1 &= 2^{-9} \int e_{k_{[1234567]}} \psi_0(k_{[12345]}, k_{[67]}) \psi_{<}(k_{[123]}, k_{[45]}) (1_{|k_{[12345]}|_\infty > N} 1_{|k_{[45]}|_\infty \leq N} + 1_{|k_{[12345]}|_\infty \leq N} 1_{N < |k_{[45]}|_\infty \leq 3N}) \\
&\quad \int_0^t \int_0^t d\sigma d\bar{\sigma} P_{t-\sigma}^\varepsilon(k_{[123]}) \Pi_{i=1}^3 P_{\sigma-s_i}^\varepsilon(k_i) p_{t-\bar{\sigma}}^\varepsilon(k_{[45]}) \varphi(\varepsilon k_{[45]}) \Pi_{i=4}^5 P_{\bar{\sigma}-s_i}^\varepsilon(k_i) \Pi_{i=6}^7 P_{t-s_i}^\varepsilon(k_i) W(d\eta_{1234567}) \\
&:= \int G(t, x, \eta_{1234567}) W(d\eta_{1234567}), \\
I_t^2 &= \sum_{i=1}^3 I_t^{2i}, I_t^{21} = 6 \int \int G(t, x, \eta_{123(-3)567}) d\eta_3 W(d\eta_{12567}), \\
I_t^{22} &= 6 \int G(t, x, \eta_{12345(-3)7}) d\eta_3 W(d\eta_{12457}), I_t^{23} = 4 \int G(t, x, \eta_{12345(-5)7}) d\eta_5 W(d\eta_{12347}), \\
I_t^3 &= \sum_{i=1}^6 I_t^{3i}, I_t^{31} = 6 \int \int G(t, x, \eta_{123(-3)(-2)67}) d\eta_{23} W(d\eta_{167}), \\
I_t^{32} &= 24 \int \int G(t, x, \eta_{123(-3)5(-2)7}) d\eta_{23} W(d\eta_{157}), I_t^{33} = 12 \int \int G(t, x, \eta_{123(-3)5(-5)7}) d\eta_{35} W(d\eta_{127}) \\
I_t^{34} &= 6 \int \int G(t, x, \eta_{12345(-2)(-3)}) d\eta_{23} W(d\eta_{145}), I_t^{35} = 12 \int \int G(t, x, \eta_{12345(-3)(-4)}) d\eta_{34} W(d\eta_{125}) \\
I_t^{36} &= 2 \int \int G(t, x, \eta_{12345(-4)(-5)}) d\eta_{45} W(d\eta_{123}), \\
I_t^4 &= \sum_{i=1}^3 I_t^{4i}, I_t^{41} = 12 \int \int G(t, x, \eta_{123(-1)(-2)(-3)7}) d\eta_{123} W(d\eta_7), \\
I_t^{42} &= 12 \int \int G(t, x, \eta_{123(-3)5(-2)(-1)}) d\eta_{123} W(d\eta_5), I_t^{43} = 24 \int \int G(t, x, \eta_{123(-3)5(-5)-2}) d\eta_{235} W(d\eta_1),
\end{aligned}$$

and J_t^1 is defined similarly as I_t^1 with $k_{[123]}, k_{[12345]}, e_{k_{[1234567]}}$ replaced by $\tilde{k}_{[123]}, \tilde{k}_{[12345]}, e_{\tilde{k}_{[1234567]}}$ respectively and $J_t^i, i = 2, 3, 4$ is defined similarly as I_t^i with the G replaced by that associated with J^1 .

Terms in the seventh chaos Now we have

$$\begin{aligned}
E|\Delta_q I_t^1|^2 &\lesssim \int \theta(2^{-q} k_{[1234567]}) \psi_0(k_{[12345]}, k_{[67]}) \psi_{<}(k_{[123]}, k_{[45]}) (1_{|k_{[12345]}|_\infty > N, |k_{[45]}|_\infty \leq N} + 1_{|k_{[12345]}|_\infty \leq N, |k_{[45]}|_\infty > N}) \\
&\quad 1_{|k_{[1234567]}| \lesssim N} \Pi_{i=1}^7 \frac{1}{|k_i|^2} \frac{1}{|k_{[123]}|^2 |k_{[45]}|^2 (|k_{[123]}|^2 + \sum_{i=1}^3 |k_i|^2) (|k_{[45]}|^2 + \sum_{i=4}^5 |k_i|^2)} dk_{1234567},
\end{aligned}$$

Observe that $|k_{45}|_\infty \simeq N$ on the support of $\psi_{<}(k_{[123]}, k_{[45]}) 1_{|k_{[12345]}|_\infty > N}$, which combining with Lemma 6.1 implies that the above term can be bounded by

$$\begin{aligned}
&\int \theta(2^{-q} k_{[1234567]}) 1_{|k_{[45]}|_\infty \simeq N, 2^q \lesssim N} \frac{1}{|k_{[123]}|^4 |k_{[45]}|^5 |k_{[67]}|} dk_{[123][45][67]} \\
&\lesssim \int 1_{2^q \lesssim N} \theta(2^{-q} k_{[1234567]}) \frac{N^{-2-\kappa}}{|k_{[12345]}|^{3-\kappa}} \frac{1}{|k_{[67]}|} dk_{[12345][67]} \lesssim \varepsilon^\kappa 2^{2q\kappa}.
\end{aligned}$$

Terms in the fifth chaos: Consider I_t^{21} first: by the formula we know that $|k_5 - k_3| \asymp N$

$$\begin{aligned}
& E|\Delta_q I_t^{21}|^2 \\
& \lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[12567]}) \Pi_{i=5}^7 \frac{1}{|k_i|^2} \frac{1}{|k_1|^2 |k_2|^2} \left[\left(\int \frac{1}{|k_3|^2 (|k_5 - k_3|^2 + |k_3|^2) (|k_{[123]}|^2 + |k_5 - k_3|^2)} dk_3 \right)^2 \right. \\
& \quad \left. + \left(\int \frac{1}{|k_3|^2 (|k_5 - k_3|^2 + |k_{[123]}|^2) (|k_{[123]}|^2 + |k_3|^2)} dk_3 \right)^2 \right] 1_{\{|k_5 - k_3| \asymp N, |k_5| \leq N\}} dk_{12567} \\
& \lesssim \int \theta(2^{-q} k_{[12567]}) 1_{2^q \lesssim N} \frac{N^{-4+2\kappa}}{|k_{[12]}|^{3-\kappa} |k_5|^{2+2\kappa} |k_{[67]}|} dk_{[12]5[67]} \lesssim \varepsilon^\kappa 2^{2q\kappa}.
\end{aligned}$$

Here in the first inequality we consider $\sigma \leq \bar{\sigma}$ and $\sigma \geq \bar{\sigma}$ separately and we used $|k_{[123]}|^2 + |k_3|^2 \gtrsim |k_{[12]}|^2$ in the second inequality.

Now we consider I_t^{22} and in this case we have $|k_{[45]}| \asymp N$, which implies that

$$\begin{aligned}
E|\Delta_q I_t^{22}|^2 & \lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[12457]}) 1_{|k_{45}| \asymp N} \frac{1}{|k_1|^2 |k_2|^2 |k_{[45]}|^5 |k_7|^2} \left(\int \frac{1}{(|k_{[123]}|^2 + |k_3|^2) |k_3|^2} dk_3 \right)^2 dk_{12[45]7} \\
& \lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[12457]}) \frac{N^{-2-\kappa}}{|k_{[12]}|^{3-\kappa} |k_{[45]}|^{3-\kappa} |k_7|^2} dk_{[12][45]7} \lesssim \varepsilon^\kappa 2^{2q\kappa}.
\end{aligned}$$

Here in the second inequality we used $|k_{[123]}|^2 + |k_3|^2 \gtrsim |k_{[12]}|^2$.

For I_t^{23} we have $|k_{[45]}| \asymp N$

$$\begin{aligned}
E|\Delta_q I_t^{23}|^2 & \lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[12347]}) 1_{\{|k_{[45]}|_\infty \asymp N\}} \Pi_{i=1}^4 \frac{1}{|k_i|^2} \frac{1}{|k_7|^2} \\
& \quad \frac{1}{(|k_{[123]}|^2 + \sum_{i=1}^3 |k_i|^2) |k_{[123]}|^2} \left(\int \frac{1}{(|k_{[45]}|^2 + |k_4|^2) |k_5|^2} dk_5 \right)^2 dk_{12347} \\
& \lesssim \int 1_{2^q \lesssim N} \theta(2^{-q} k_{[12347]}) \frac{1}{|k_{[1234]}|^{2-\kappa}} \frac{N^{-2+\kappa}}{|k_7|^2} dk_{[1234]7} \lesssim \varepsilon^\kappa 2^{2q\kappa},
\end{aligned}$$

where we used $|k_{[45]}|^2 + |k_4|^2 \gtrsim |k_5|^2$ in the second inequality.

Terms in the third chaos: For I_t^{31} , we have

$$\begin{aligned}
E|\Delta_q I_t^{31}|^2 & \lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[167]}) \psi_0(k_1, k_{[67]}) \psi_{<}(k_{[123]}, k_{[23]}) (1_{|k_1|_\infty > N, |k_{[23]}|_\infty \leq N} + 1_{|k_1|_\infty \leq N, N < |k_{[23]}|_\infty \leq 3N}) \\
& \quad \frac{1}{|k_1|^2} \frac{1}{|k_6|^2 |k_7|^2} \left(\int \frac{1}{(|k_{[123]}|^2 + |k_{[23]}|^2) |k_{[23]}|^2 |k_2|^2 |k_3|^2} dk_{[23]} \right)^2 dk_{167} \\
& \lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[167]}) \frac{N^{-3-\kappa}}{|k_1|^{3-\kappa} |k_{[67]}|} dk_{1[67]} \lesssim \varepsilon^\kappa 2^{2q\kappa}.
\end{aligned}$$

Here we used $|k_{[23]}|^2 \lesssim |k_2|^2 + |k_3|^2$ in the first inequality and in the second inequality we used $|k_{[23]}| \asymp N$.

For I_t^{32} we have

$$\begin{aligned}
E|\Delta_q I_t^{32}|^2 &\lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[157]}) 1_{\{|k_5 - k_3|_\infty \asymp N\}} \frac{1}{|k_1|^2 |k_5|^2 |k_7|^2} 1_{|k_5| \leq N} \\
&\quad \left[\left(\int \frac{1}{(|k_5 - k_3|^2 + |k_3|^2)(|k_{[123]}|^2 + |k_2|^2 + |k_5 - k_3|^2) |k_2|^2 |k_3|^2} dk_{23} \right)^2 \right. \\
&\quad \left. + \left(\int \frac{1}{(|k_{[123]}|^2 + |k_2|^2 + |k_3|^2)(|k_{[123]}|^2 + |k_2|^2 + |k_5 - k_3|^2) |k_2|^2 |k_3|^2} dk_{23} \right)^2 \right] dk_{157} \\
&\lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[157]}) N^{-3} \frac{1}{|k_1|^2 |k_5|^{3-\kappa} |k_7|^2} dk_{157} \lesssim \varepsilon^\kappa 2^{2q\kappa},
\end{aligned}$$

Here in the first inequality we consider $\sigma \leq \bar{\sigma}$ and $\sigma \geq \bar{\sigma}$ separately. For I_t^{33} we have

$$\begin{aligned}
E|\Delta_q I_t^{33}|^2 &\lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[127]}) 1_{|k_5 - k_3| \asymp N, |k_{[12]}| \lesssim N} \frac{1}{|k_1|^2 |k_2|^2 |k_7|^2} \\
&\quad \left[\left(\int \frac{1}{(|k_3|^2 + |k_5 - k_3|^2 + |k_5|^2)(|k_{[123]}|^2 + |k_5 - k_3|^2 + |k_5|^2) |k_3|^2 |k_5|^2} dk_{35} \right)^2 \right. \\
&\quad \left. + \left(\int \frac{1}{(|k_3|^2 + |k_{[123]}|^2)(|k_{[123]}|^2 + |k_5 - k_3|^2 + |k_5|^2) |k_3|^2 |k_5|^2} dk_{35} \right)^2 \right] dk_{127} \\
&\lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[127]}) \frac{N^{-2+2\kappa}}{|k_{[12]}|^{3-\kappa} |k_7|^2} dk_{[12]7} \lesssim \varepsilon^\kappa 2^{2q\kappa}.
\end{aligned}$$

For I_t^{34} we have

$$\begin{aligned}
E|\Delta_q I_t^{34}|^2 &\lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[145]}) 1_{|k_{[45]}| \asymp N} \frac{1}{|k_1|^2 |k_{[45]}|^5} \\
&\quad \left(\int \frac{1_{|k_{[23]}| \lesssim N}}{(|k_{[123]}|^2 + \sum_{i=2}^3 |k_i|^2) |k_2|^2 |k_3|^2} dk_{23} \right)^2 dk_{145} \\
&\lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[145]}) 1_{|k_{[45]}|_\infty \asymp N} \frac{N^\kappa}{|k_{[45]}|^5 |k_1|^2} dk_{1[45]} \lesssim \varepsilon^\kappa 2^{2q\kappa}.
\end{aligned}$$

For I_t^{35} we have

$$\begin{aligned}
E|\Delta_q I_t^{35}|^2 &\lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[125]}) 1_{|k_{[45]}|_\infty \asymp N} \frac{1}{|k_1|^2 |k_2|^2 |k_5|^2} \\
&\quad \left(\int \frac{1}{(|k_{[45]}|^2 + |k_4|^2)(|k_{[123]}|^2 + |k_3|^2) |k_3|^2 |k_4|^2} dk_{34} \right)^2 dk_{125} \\
&\lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[125]}) \frac{N^{-2+\kappa}}{|k_5|^2 |k_{[12]}|^{3-\kappa}} dk_{[12]5} \lesssim \varepsilon^\kappa 2^{2q\kappa}.
\end{aligned}$$

Here in the second inequality we used $|k_{[123]}|^2 + |k_3|^2 \gtrsim |k_{[12]}|^2$.

For I_t^{36} we have

$$\begin{aligned}
E|\Delta_q I_t^{36}|^2 &\lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[123]}) \frac{1}{|k_{[123]}|^4} \Pi_{i=1}^3 \frac{1}{|k_i|^2} \\
&\quad \left(\int \frac{(1_{|k_{[12345]}|_\infty > N} 1_{|k_{[45]}|_\infty \leq N} + 1_{|k_{[12345]}|_\infty \leq N} 1_{N < |k_{[45]}|_\infty \leq 3N})}{(|k_{[45]}|^2 + \sum_{i=4}^5 |k_i|^2) |k_4|^2 |k_5|^2} dk_{45} \right)^2 dk_{123},
\end{aligned}$$

Now we use similar argument as Section 6.3. In the first case without loss of generality we assume that $|k_{[123]}^i + k_{[45]}^i| > N$ for some i . Then there are at most $|k_{[123]}^i|$ values of $k_{[45]}^i$ with $|k_{[12345]}^i| > N$ and $|k_{[45]}^i| \leq N$. In the second case similarly we obtain that there are at most $|k_{[123]}^i|$ values of $k_{[45]}^i$ with $|k_{[45]}^i| > N$ and $|k_{[12345]}^i| \leq N$. Thus we obtain

$$E|\Delta_q I_t^{36}|^2 \lesssim 1_{2^q \lesssim N} \int \theta(2^{-q} k_{[123]}) N^{-2+\kappa} \frac{1}{|k_{[123]}|^2} dk_{[123]} \lesssim \varepsilon^\kappa 2^{q\kappa}.$$

Terms in the first chaos: For I^{41} we obtain that

$$\begin{aligned} E|\Delta_q I_t^{41}|^2 &\lesssim \int \theta(2^{-q} k_7) \frac{1}{|k_7|^2} \left[\int \frac{1_{|k_{[12]}| \leq N}}{|k_2|^2 |k_3|^2 |k_1|^2 (|k_{[123]}|^2 + |k_3|^2) |k_{[12]}|^2} dk_{123} \right]^2 dk_7 \\ &\lesssim \int \theta(2^{-q} k_7) \frac{N^{-2+\kappa}}{|k_7|^2} dk_7 \lesssim \varepsilon^\kappa 2^{2q\kappa}. \end{aligned}$$

For I^{42} we obtain that

$$\begin{aligned} E|\Delta_q I_t^{42}|^2 &\lesssim \int \theta(2^{-q} k_5) \left(\int 1_{|k_5 - k_3| \leq N} \Pi_{i=1}^3 \frac{1}{|k_i|^2} \frac{1_{\{|k_i| \leq N, i=1,2,3\}}}{(|k_{[123]}|^2 + \sum_{i=1}^2 |k_i|^2) (|k_{[5-3]}|^2 + |k_3|^2)} dk_{123} \right)^2 dk_5 \\ &\lesssim \int \theta(2^{-q} k_5) \frac{N^{-2+\kappa}}{|k_5|^2} dk_5 \lesssim \varepsilon^\kappa 2^{2q\kappa}. \end{aligned}$$

For I^{43} we obtain that

$$\begin{aligned} E|\Delta_q I_t^{43}|^2 &\lesssim \int \theta(2^{-q} k_1) \frac{1}{|k_1|^2} \left(\int \frac{1_{|k_{5-3}| \leq N}}{|k_2|^2 |k_3|^2 |k_5|^2 (|k_{[123]}|^2 + |k_2|^2 + |k_3|^2) (|k_{[5-3]}|^2 + |k_5|^2)} dk_{235} \right)^2 dk_1 \\ &\lesssim \int \theta(2^{-q} k_1) \frac{N^{-2+\kappa}}{|k_1|^2} dk_1 \lesssim \varepsilon^\kappa 2^{2q\kappa}. \end{aligned}$$

Moreover for J_t^i and other terms in D^N we could use similar calculations and Lemma 6.5 to obtain the same estimates. Then by using Gaussian hypercontractivity, Lemma 2.1 and Kolomogorov continuity criterion we obtain that $D_N \xrightarrow{P} 0$ as $\varepsilon \rightarrow 0$.

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